Quantum dynamics of a wave packet for a Cantor spectrum

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMANS¹

1 - Technion Institute of Technology ; 2 - École Normale Supérieure Cachan

IPS meeting 2014 - Be'er Sheva

Introduction

- 2 Analytical study in the case of a Cantor set
 - Appropriate measure
 - Time-averaged return probability
 - Other relevant quantities

3 Experiment : cavity polaritons in a Fibonacci potential

4 Summary

Analytical study in the case of a Cantor set Experiment : cavity polaritons in a Fibonacci potential Summary

Table of contents



2 Analytical study in the case of a Cantor set

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

・ 同 ト ・ ヨ ト ・ ヨ

Analytical study in the case of a Cantor set Experiment : cavity polaritons in a Fibonacci potential Summary

Table of contents



2 Analytical study in the case of a Cantor set

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

・ 同 ト ・ ヨ ト ・ ヨ

Analytical study in the case of a Cantor set Experiment : cavity polaritons in a Fibonacci potential Summary

Table of contents



- 2 Analytical study in the case of a Cantor set
- 3 Experiment : cavity polaritons in a Fibonacci potential

4 Summary

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

4 冊 ト 4 三 ト 4 三 ト

Consider systems with a (spatial and/or spectral) fractal feature.

Motivation :

- singular continuous spectra are ubiquitous in physical systems, but not well known
- Explain experimental and numerical observations

Objectives :

- Identify typical behaviour of relevant physical quantities in the presence of a fractal feature (spatial and/or spectral) \rightarrow power-law modulated by log-periodic oscillations
- Derive analytical expressions, and relate them to experimental results

4 日 2 4 周 2 4 月 1 4 月 1

Analytical study in the case of a Cantor set Experiment : cavity polaritons in a Fibonacci potential Summary

Fractal :

- a set with a discrete scaling symmetry
- fractal dimension d_f (length scale l; volume scales as l^{d_f})

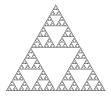


Figure: Sierpinski gasket ; $d_f = \frac{\ln 3}{\ln 2}$

- A - E - N

Table of contents

Appropriate measure Time-averaged return probability Other relevant quantities

(1日) (1日) (1日)



2 Analytical study in the case of a Cantor set

• Appropriate measure

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

Table of contents

Appropriate measure Fime-averaged return probability Other relevant quantities

A (1) < A (1) < A (1) < A (1) </p>

1 Introduction

2 Analytical study in the case of a Cantor set

- Appropriate measure
- Time-averaged return probability

Table of contents

Appropriate measure Time-averaged return probability Other relevant quantities

A (1) < A (1) < A (1) < A (1) </p>

1 Introduction

2 Analytical study in the case of a Cantor set

- Appropriate measure
- Time-averaged return probability
- Other relevant quantities

3 Experiment : cavity polaritons in a Fibonacci potential

④ Summary

Quantum dynamics for a triadic Cantor set spectrum.

 $H = -\frac{d^2}{dx^2} + V(x)$ on $L^2(\mathbb{R})$, with associated triadic Cantor set spectrum.



Figure: Triadic Cantor set. Fractal dimension $d_f = \frac{\ln 2}{\ln 3}$

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

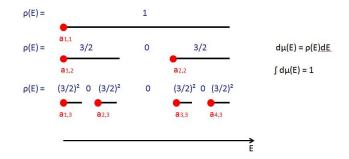
イロト イポト イヨト イヨト

-

Appropriate measure Time-averaged return probability Other relevant quantities

イロト イポト イヨト イヨト

Triadic Cantor set K : zero (Lebesgue) measure. Degenerate density of states $\rho(\epsilon)$. Define measure : $d\mu(\epsilon) \equiv \rho(\epsilon)d\epsilon$:



Verifies : $\int_{K} f(x) d\mu(x) = \frac{1}{2} \int_{K} f(\frac{x}{3}) d\mu(x) + \frac{1}{2} \int_{K} f(\frac{x+2}{3}) d\mu(x)$ (scaling property)

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

マロト イヨト イヨト

How does the fractal feature affect the dynamics of the wave packet $|\psi(t)
angle$?

- RMS displacement $\Delta x(t) = \sqrt{\int (x(t) x_0)^2 |\psi(x, t)|^2 dx}$
- 2 return probability $p(t) = |\langle \psi(0) | \psi(t) \rangle|^2$
- 3 ...

A relevant quantity : the time-averaged return probability :

$$C(t) = rac{1}{t}\int_0^t |\langle \psi(0)|\psi(t')
angle|^2 dt'$$

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrum

(日) (同) (三) (三)

A relevant quantity : the time-averaged return probability :

$$C(t) = \frac{1}{t} \int_0^t |\langle \psi(0) | \psi(t') \rangle|^2 dt' = \iint_{K \times K} \operatorname{sinc}((\epsilon' - \epsilon)t) d\mu(\epsilon) d\mu(\epsilon')$$

Note : ϵ and t are coupled variables. Method

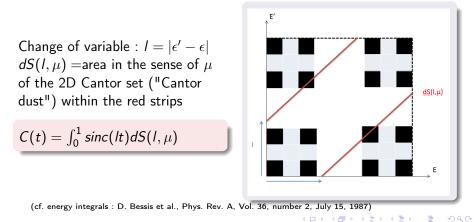
- use the scaling symmetry of the measure to obtain a functional equation
- 2 apply the Mellin transform

• • = • • = •

Introduction Analytical study in the case of a Cantor set

Experiment : cavity polaritons in a Fibonacci potential Summary Appropriate measure Time-averaged return probability Other relevant quantities

$$C(t) = \iint_{K imes K} sinc((\epsilon' - \epsilon)t) d\mu(\epsilon) d\mu(\epsilon')$$



Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrum

Appropriate measure Time-averaged return probability Other relevant quantities

The main steps :

- **9** Mellin transform : $\zeta(\alpha) = \int_0^1 I^{-\alpha} dS(I,\mu) = \iint_{K \times K} \frac{d\mu(x)d\mu(y)}{|x-y|^{\alpha}}$
- Scaling property : $\int_{K \times K} \frac{d\mu(x)d\mu(y)}{|x-y|^{\alpha}} = \frac{1}{2} \iint_{K \times K} \frac{d\mu(x)d\mu(y)}{3^{-\alpha}|x-y|^{\alpha}} + \underbrace{\frac{1}{2} \iint_{K \times K} \frac{d\mu(x)d\mu(y)}{3^{-\alpha}|x-y+2|^{\alpha}}}_{g(\alpha)}$

$$=\frac{3^{\alpha}}{2}\zeta(\alpha)+g(\alpha)$$

- - Mellin transform of $f: M_f(s) = \int_0^\infty f(x) x^{s-1} dx$
 - Inverse Mellin transform : $M_{\phi}^{-1}(x) = rac{1}{2i\pi}\int_{c-i\infty}^{c+i\infty}\phi(s)x^{-s}ds$

Appropriate measure Time-averaged return probability Other relevant quantities

Time-averaged return probability for triadic Cantor set

$$C(t) = \frac{t^{-\ln 2/\ln 3}}{\ln(3)} \sum_{n} \epsilon_n e^{-\frac{2i\pi n}{\ln(3)}\ln t} \sin\left(\frac{\pi}{2}(d_n - 1)\right) \Gamma(d_n - 1)$$

$$\begin{aligned} d_n &= \frac{\ln(2)}{\ln(3)} + \frac{2i\pi n}{\ln(3)}, \ \epsilon_n = \int_{K \times K} \frac{d\mu(x)d\mu(y)}{|x-y+2|^{d+\frac{2i\pi n}{\ln(3)}}}, \ \Gamma \ \text{the Euler gamma function} \\ \Rightarrow C(t) &= t^{-\ln 2/\ln 3} G_p(\ln t/\ln 3), \ G_p(y+1) = G_p(y) \end{aligned}$$

power-law modulated by log-periodic oscillations.

[●] C(t) ∝ t^{-d_f}, with d_f = $\frac{\ln 2}{\ln 3}$ (fractal dimension) for a triadic Cantor set.

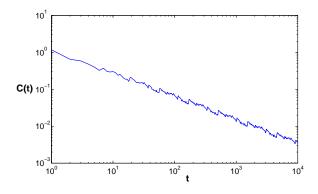


Figure: Return probability C(t) for a Cantor spectrum

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

Introduction Analytical study in the case of a Cantor set

Experiment : cavity polaritons in a Fibonacci potential Summary Appropriate measure Time-averaged return probability Other relevant quantities

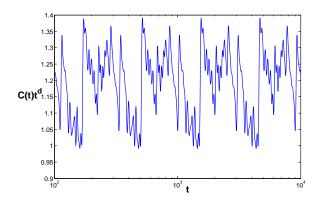


Figure: Return probability divided by $t^{-\ln 2/\ln 3}$, $C(t)t^{\ln 2/\ln 3}$ for a Cantor spectrum

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

Appropriate measure Time-averaged return probability Other relevant quantities

・ 同 ト ・ ヨ ト ・ ヨ ト

Partial summary :

• General method :

scaling property

- \rightarrow functional equation $f(x) = \frac{1}{b}f(ax) + g(x)$
- \rightarrow Mellin transform

$$ightarrow f(x) = x^{\ln(b)/\ln(a)} G_{
ho}\left(rac{\ln x}{\ln a}
ight), \ G_{
ho}(y+1) = G_{
ho}(y)$$

- Typical behaviour due to a fractal spectrum : power-law modulated by log-periodic oscillations
- The method gives an intuition of which are the relevant physical quantities to observe typical behaviour

Other relevant quantities :

- Participation ratio : $P_r(t) = \frac{1}{t} \int_0^t \frac{dt}{|\langle x|\psi(t')\rangle|^4}$
- RMS displacement : $\Delta x(t) = \sqrt{\int (x(t) x_0)^2 |\psi(x, t)|^2 dx}$

・ 同 ト ・ ヨ ト ・ ヨ ト

Table of contents



2 Analytical study in the case of a Cantor set

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

▲ □ ▶ ▲ □ ▶ ▲ □ ▶

Table of contents



2 Analytical study in the case of a Cantor set

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

▲ □ ▶ ▲ □ ▶ ▲ □ ▶

Table of contents



- 2 Analytical study in the case of a Cantor set
- Experiment : cavity polaritons in a Fibonacci potential

4 Summary

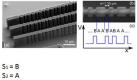
Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

• • • • • • • • •

Cavity polaritons (quasi-particles arising from the strong coupling between photons and excitons) placed in a guasi-periodic potential.

Mapping to a 1D problem with effective quasi-periodic potential V(x) with a Cantor like spectrum :

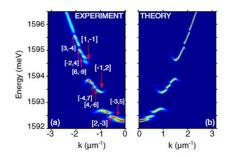
$$E\psi(x) = \frac{\hbar^2}{2m_{ph}} \left[-\frac{d^2}{dx^2} + V(x) \right] \psi(x)$$



S3 = BA	$S_{i+1} = S_{i-1}S_i$
$S_4 = ABA$	-)))
S5 = BAABA	
S6 = ABABAABA	

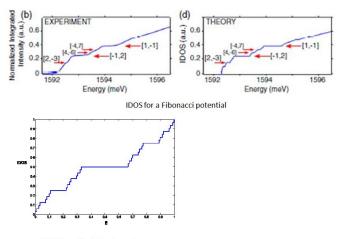
Figure: Fibonacci cavity

Numerical and experimental results



Gaps of the energy spectrum for a Fibonacci potential (gap-labeling theorem)

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur



IDOS for a triadic Cantor set

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrum

→ < Ξ →</p>

< 口 > < 同

< ∃ >

э

Time-averaged return probability : $C(t) = \frac{1}{t} \int_0^t |\langle \psi(0) | \psi(t') \rangle|^2 dt'$

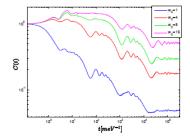


Figure: Time-averaged return probability of a wave packet in a Fibonacci potential.

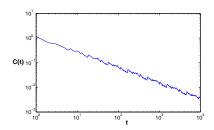


Figure: Time-averaged return probability for a triadic Cantor spectrum.

A 3 b

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

Table of contents



2 Analytical study in the case of a Cantor set

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

▲ □ ▶ ▲ □ ▶ ▲ □ ▶

Table of contents



2 Analytical study in the case of a Cantor set

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

▲ □ ▶ ▲ □ ▶ ▲ □ ▶

Table of contents

1 Introduction

2 Analytical study in the case of a Cantor set

3 Experiment : cavity polaritons in a Fibonacci potential

4 Summary

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

・ 同 ト ・ ヨ ト ・ ヨ

In summary :

- We identified the signature of a fractal spectrum through measurable physical quantities : power-law modulated by log-periodic oscillation, for dynamical and spectral quantities (C(t), integrated density of states)
- General method to study theoretically systems with a fractal feature
- Oscillations difficult to observe ; easier with dynamical quantities
- Verified experimentally with polaritons in a Fibonacci potential (Cantor like spectrum)

・ 同 ト ・ ヨ ト ・ ヨ ト

Next : other relevant quantities in quantum mechanics

- Signature of the fractal spectral feature in the Thouless coefficient $\gamma(\epsilon)$
- Spatial correlation of the eigenfunction; generalization of Thouless formula : $\overline{\phi(0,\epsilon)^*\phi(x,\epsilon)} \propto e^{-|x|\gamma(\epsilon)}$
- Spontaneous emission of a single atom coupled to a fractal spectrum
- Superradiance of a two atoms system coupled to a fractal spectrum

・ 同 ト ・ ヨ ト ・ ヨ ト …

Questions?

Ariane SORET^{1,2} Evgeni GUREVICH¹, Eric AKKERMAN Quantum dynamics of a wave packet for a Cantor spectrur

< ロ > < 同 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ >

э