Introduction and First Notations Polar divisor considerations Proof of the Main Theorem Acknowledgements

# Hidden Functional Equations for Rankin-Selberg transforms: New Results

David LECOMTE

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# Outline

- Introduction and First Notations
  - The SL<sub>2</sub>(Z) Eisenstein Series
  - The Hilbert Modular Eisenstein Series
  - Zagier's Renormalization
  - Previous Results and Present Goal
- 2 Polar divisor considerations
  - Identifying the Polar Divisor of R<sub>A</sub>
  - Finding the Group of Π
  - An Upper Bound for |S|
- Proof of the Main Theorem
  - More Notations
  - First Computations
  - Along Came Poisson
  - Wrap It Up, David



Polar divisor considerations Proof of the Main Theorem Acknowledgements The SL<sub>2</sub>(ℤ) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

# The $SL_2(\mathbb{Z})$ Eisenstein Series

• 
$$E^*(z,s) = \frac{1}{2}\pi^{-s}\Gamma(s)\sum_{\substack{(m,n)\in\mathbb{Z}^2\\(m,n)\neq 0}}\frac{y^s}{|mz+n|^{2s}}$$
  $z\in\mathscr{H}$  Ress

•  $E^*(z, s)$  is  $SL_2(\mathbb{Z})$ -automorphic in the variable *z*.

•  $E^*(z, s)$  is analytic in the variable s.

Polar divisor considerations Proof of the Main Theorem Acknowledgements The SL<sub>2</sub>(Z) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The SL<sub>2</sub>(ℤ) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The SL<sub>2</sub>(ℤ) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The SL<sub>2</sub>(Z) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

## **Fourier Expansion**

$$E^{*}(z,s) = y^{s}\zeta^{*}(2s) + y^{1-s}\zeta^{*}(2-2s) + 2\sqrt{y}\sum_{n\neq 0}\tau_{s-\frac{1}{2}}(|n|)K_{s-\frac{1}{2}}(2\pi|n|y)e^{2\pi inx}$$

where

• 
$$\tau_{\omega}(n) = n^{-\omega} \sum_{\substack{d \mid n \\ d > 0}} d^{2\omega}$$
 is a divisor sum;  
•  $\mathcal{K}_{\omega}(y) = \frac{1}{2} \int_{0}^{+\infty} e^{-\frac{y}{2}(t+\frac{1}{t})} t^{\omega} \frac{dt}{t}$  is a Bessel function.

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The SL<sub>2</sub>(Z) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The SL<sub>2</sub>(Z) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The SL<sub>2</sub>(Z) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The SL<sub>2</sub>(Z) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

### What You Need to Know About $\tau_{\omega}$ and $K_{\omega}$

• 
$$\tau_{\omega}(n) = \tau_{-\omega}(n)$$
 and  $K_{\omega}(y) = K_{-\omega}(y)$ ;

•  $K_{\omega}$  has rapid decay:

$$\forall y > 2$$
  $|K_{\omega}(y)| \leq M e^{-\frac{y}{2}}$ 

As a consequence,  $E^*(z, s)$  is the uniform limit of its Fourier series, can be analytically continued to  $\mathbb{C} \setminus \{0, 1\}$  and satisfies the functional equation

$$\forall z \in \mathscr{H} \quad \forall s \neq 0, 1 \qquad E^*(z, s) = E^*(z, 1-s)$$

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The SL<sub>2</sub>(ℤ) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Introduction and First Notations Polar divisor considerations Proof of the Main Theorem The SL<sub>2</sub>(Z) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

# The Setting

- Q → K is a totally real field extension with ring of integers 
   *o* and discriminant D.
- There are exactly *N* distinct embeddings  $K \hookrightarrow \mathbb{C}$ , denoted as  $\sigma_1, \ldots, \sigma_N$ .
- For convenience, we write also

$$\forall \alpha \in \mathbf{K} \qquad \alpha^{(i)} = \sigma_i(\alpha)$$

• A is an ideal class in K and  $\alpha$  an ideal in  $A^{-1}$ .

Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements

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The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

### The Hilbert Modular Eisenstein Series

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$$E_{K,A}^{\star}(z,s) = \mathbb{N}(\mathfrak{a})^{2s} \pi^{-Ns} \Gamma(s)^{N} D^{s} \sum_{\substack{(\alpha,\beta) \in \mathfrak{a}^{2}/\mathfrak{o}^{\times} \\ (\alpha,\beta) \neq 0}} \prod_{i=1}^{N} \frac{Y_{i}^{s}}{\left|\alpha^{(i)} z + \beta^{(i)}\right|^{2s}}$$

where  $z \in \mathscr{H}$  and Ress > 1.

- $E_{K,A}^{\star}(z,s)$  is  $SL_2(\mathbb{Z})$ -automorphic in the variable z;
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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

### **Fourier Expansion**

$$\begin{split} E_{K,A}^{\star}(z,s) &= y^{Ns} \zeta_{K,A}^{\star}(2s) + y^{N(1-s)} \zeta_{K,A}^{\star}(2-2s) \\ &+ 2^{N} y^{\frac{N}{2}} \sum_{\substack{\xi \in \mathfrak{D}^{-1} \\ \xi \neq 0}} \tau_{s-\frac{1}{2}}^{K,A}(\xi \mathfrak{D}) \Big(\prod_{i=1}^{N} K_{s-\frac{1}{2}}(2\pi y |\xi^{(i)}|) \Big) e^{2\pi i x \operatorname{Tr} \xi} \end{split}$$

where  $\tau_{\omega}(\mathfrak{c}) = \mathbb{N}(\mathfrak{c})^{-\omega} \sum_{\substack{\mathfrak{b} \text{ ideal in A} \\ \mathfrak{b}|\mathfrak{c}}} \mathbb{N}(\mathfrak{b})^{2\omega}$  is a generalized divisor sum.

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

### Zagier's Theorem

#### Theorem

Let F be an SL<sub>2</sub>( $\mathbb{Z}$ )-automorphic function on  $\mathcal{H}$ . Suppose that there exists a function  $\varphi$  of the form

$$arphi(\mathbf{y}) = \sum_{i=1}^{\ell} c_i \mathbf{y}^{lpha_i} \ln^{n_i} \mathbf{y} \qquad lpha_i \in \mathbb{C} \quad n_i \in \mathbb{N}$$

such that  $F(z) - \varphi(y) = \underset{y \to +\infty}{=} o(y^{-n})$  for every positive integer n. Then

$$R(s) = RN \int_{SL_2(\mathbb{Z})\setminus\mathscr{H}} E^*(z,s)F(z) \frac{dz}{y^2} = \int_{0}^{+\infty} (a_0(y) - \varphi(y))y^{s-1} \frac{dy}{y}$$

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

# Zagier's Theorem: Ze Continuation

#### Theorem

where 
$$a_0(y) = \int_0^1 F(x + iy) dx$$
, is well defined as an absolutely convergent integral for Res big enough. It has analytic continuation to  $\mathbb{C}$ , except for poles at 0, 1, the  $\alpha_i$ 's and the  $(1 - \alpha_i)$ 's. Furthermore,  $R(s) = R(1 - s)$ .

Polar divisor considerations Proof of the Main Theorem Acknowledgements The SL<sub>2</sub>(Z) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

## **Previous Results**

#### Were already studied by Professor BUMP:

•  $RN \int_{SL_2(\mathbb{Z})\setminus\mathscr{H}} E^*(z, s_0) E^*_{K,A}(z, s_1) \frac{dz}{y^2}$  when *K* is a totally real cubic field and *A* is an ideal class;

• 
$$RN \int_{SL_2(\mathbb{Z})\setminus\mathcal{H}} E^*(z,s_0)E^*(z,s_1)E^*(z,s_2)E^*(z,s_3)\frac{\mathrm{d}z}{y^2}$$

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

# **Our Goal**

We suspect that a similar phenomenon occurs and want to identify the full group of functional equations for the renormalized Rankin-Selberg transform

$$R_A(s_0, s_1, s_2) = \int_{SL_2(\mathbb{Z})\setminus \mathscr{H}} E^*(z, s_0) E^*(z, s_1) E^*_{K,A}(z, s_2) \frac{\mathrm{d}z}{y^2}$$

when K is a real quadratic field and A is an ideal class. Zagier's theorem, together with our knowledge of the objects involved, predict 16 functional equations. In fact,

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Polar divisor considerations Proof of the Main Theorem Acknowledgements The SL<sub>2</sub>(Z) Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

# **Our Goal**

We suspect that a similar phenomenon occurs and want to identify the full group of functional equations for the renormalized Rankin-Selberg transform

$$R_A(s_0, s_1, s_2) = \int\limits_{SL_2(\mathbb{Z})\backslash \mathscr{H}} E^{\star}(z, s_0) E^{\star}(z, s_1) E^{\star}_{K,A}(z, s_2) \frac{dz}{y^2}$$

when K is a real quadratic field and A is an ideal class. Zagier's theorem, together with our knowledge of the objects involved, predict 16 functional equations. In fact,

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Introduction and First Notations

Polar divisor considerations Proof of the Main Theorem Acknowledgements The  $SL_2(\mathbb{Z})$  Eisenstein Series The Hilbert Modular Eisenstein Series Zagier's Renormalization Previous Results and Present Goal

#### Main Theorem

#### Theorem

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Let K be a real quadratic field and A an ideal class. The function  $R_A(s_0, s_1, s_2)$  has a group of functional equations of order 48. It is generated by the transformations

$$s_{0} \longmapsto 1 - s_{0} \qquad s_{1} \longmapsto 1 - s_{1} \qquad s_{2} \longmapsto 1 - s_{2} \qquad s_{0} \longleftrightarrow s_{1}$$

$$nd \qquad \begin{bmatrix} s_{0} \\ s_{1} \\ s_{2} \end{bmatrix} \stackrel{w}{\mapsto} \begin{bmatrix} -\frac{s_{0}}{2} + \frac{s_{1}}{2} + s_{2} \\ 1 - \frac{s_{0}}{2} + \frac{s_{1}}{2} - s_{2} \\ 1 - \frac{s_{0}}{2} - \frac{s_{1}}{2} \end{bmatrix}.$$

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Identifying the Polar Divisor of  $R_A$ Finding the Group of  $\Pi$ An Upper Bound for |S|

### Strategy

The polar divisor  $\Pi$  of  $R_A$  is the subset of  $\mathbb{C}^3$  at which  $R_A$  is undefined. Zagier's theorem provides us with a complete description of  $\Pi$ . All we have to do is identify the part  $\varphi$  of F that is not of rapid decay and hope it looks like

$$arphi(\mathbf{y}) = \sum_{i=1}^{\ell} c_i \mathbf{y}^{lpha_i} \ln \mathbf{y}^{n_i}$$

Once we know the  $\alpha_i$ 's, we know  $\Pi$ .

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#### Reminder

#### Remember that

$$\begin{split} E^{*}(z,s_{1}) &= y^{s_{1}}\zeta^{*}(2s_{1}) + y^{1-s_{1}}\zeta^{*}(2-2s_{1}) \\ &+ 2\sqrt{y}\sum_{n\neq 0}\tau_{s_{1}-\frac{1}{2}}(|n|)K_{s_{1}-\frac{1}{2}}(2\pi|n|y)e^{2\pi inx} \\ E^{*}_{K,A}(z,s_{2}) &= y^{2s_{2}}\zeta^{*}_{K,A}(2s_{2}) + y^{2(1-s_{2})}\zeta^{*}_{K,A}(2-2s_{2}) \\ &+ 4y\sum_{\substack{\xi \in \mathfrak{D}^{-1}\\\xi \neq 0}}\tau^{K,A}_{s_{2}-\frac{1}{2}}(\xi\mathfrak{D})\Big(\prod_{l=1}^{2}K_{s_{2}-\frac{1}{2}}(2\pi y|\xi^{(l)}|)\Big)e^{2\pi ix\operatorname{Tr}\xi} \\ F(z) &= E^{*}(z,s_{1})E^{*}_{K,A}(z,s_{2}) \end{split}$$

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#### Expression of $\varphi$

## $\begin{aligned} \varphi(\mathbf{y}) &= \mathbf{y}^{s_1 + 2s_2} \zeta^*(2s_1) \zeta^*_{K,A}(2s_2) + \mathbf{y}^{2 + s_1 - 2s_2} \zeta^*(2s_1) \zeta^*_{K,A}(2 - 2s_2) \\ &+ \mathbf{y}^{1 - s_1 + 2s_2} \zeta^*(2 - 2s_1) \zeta^*_{K,A}(2s_2) + \mathbf{y}^{3 - s_1 - 2s_2} \zeta^*(2 - 2s_1) \zeta^*_{K,A}(2 - 2s_2) \end{aligned}$

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#### Cartesian Equation of Π

By Zagier's Theorem, the polar divisor of  $R_A$  is the union of the 14 hyperplanes

 $s_0 = s_1 + 2s_2$   $s_0 = 2 + s_1 - 2s_2$   $s_0 = 1 - s_1 + 2s_2$   $s_0 = 3 - s_1 - 2s_2$ 

 $s_0 = 1 - s_1 - 2s_2 \quad s_0 = -1 - s_1 + 2s_2 \quad s_0 = s_1 - 2s_2 \quad s_0 = -2 + s_1 + 2s_2$ 

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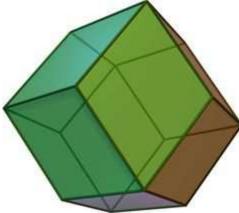
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#### So what does ∏ look like?

The 14 hyperplanes in  $\mathbb{C}^3$  cut out a rhombic dodecahedron:



Identifying the Polar Divisor of  $R_A$ Finding the Group of  $\Pi$ An Upper Bound for |S|

## The Group of Π

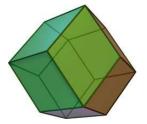
The group  $S_{\Pi}$  of transformations leaving  $\Pi$  invariant is computed through standard counting arguments.

For example,  $S_{\Pi}$  acts transitively on the set of vertices that are connected to 4 other vertices.

There are 6 such vertices.

Each of them is fixed by 8 transformations: 4 reflexions and 4 rotations.

So  $|S_{\Pi}| = 6 \times 8 = 48$ .



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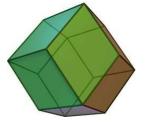
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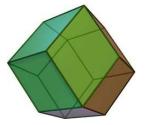
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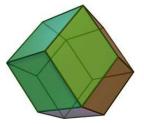
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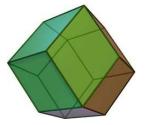
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#### Relationship Between S and $S_{\Pi}$

There is an easy inclusion between *S* and *S*<sub>Π</sub>. Every functional equation for  $R_A$  should leave  $\Pi$  invariant. So  $S \subset S_{\Pi}$  and it follows that  $|S| \leq 48$ .

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More Notations First Computations Along Came Poissor Wrap It Up, David

### The $SL_3(\mathbb{Z})$ Eisenstein series I

- $\nu_1$  and  $\nu_2$  are complex numbers with real part bigger than  $\frac{2}{3}$ .
- $\Omega$  is the set of quadruples  $(\alpha, \beta, \xi, \eta) \in \mathfrak{a} \times (\mathfrak{a}\mathfrak{D})^{-1} \times \mathbb{Z} \times \mathbb{Z}$  with
  - $\alpha$  and  $\xi$  not simultaneously 0;
  - $\beta$  and  $\eta$  not simultaneously 0;
  - Tr  $\alpha\beta + \xi\eta = 0$ .
- $c = \mathbb{N}(\mathfrak{a})^{-\frac{1}{2}} D^{-\frac{1}{4}};$
- $t_1$  and  $t_2$  are positive real numbers.

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More Notations First Computations Along Came Poisson Wrap It Up, David

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More Notations First Computations Along Came Poisson Wrap It Up, David

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#### More Notations

First Computations Along Came Poisson Wrap It Up, David

## The $SL_3(\mathbb{Z})$ Eisenstein Series II

Let 
$$G_{\nu_1,\nu_2}(t_1, t_2) = \frac{1}{4} \pi^{-\frac{3\nu_1}{2}} \Gamma\left(\frac{3\nu_1}{2}\right) \pi^{-\frac{3\nu_2}{2}} \Gamma\left(\frac{3\nu_2}{2}\right) \zeta^*(3\nu_1 + 3\nu_2 + 1)$$
  
  $\times \sum_{(\alpha,\beta,\xi,\eta)\in\Omega} \left[ \left(Ct_1\alpha^{(1)}\right)^2 + \left(Ct_2\alpha^{(2)}\right)^2 + \left(\frac{\xi}{t_1t_2}\right) \right]^{-\frac{3\nu_1}{2}} \times \left[ \left(\frac{\beta^{(1)}}{ct_1}\right)^2 + \left(\frac{\beta^{(2)}}{ct_2}\right)^2 + (\eta t_1 t_2)^2 \right]^{-\frac{3\nu_2}{2}}$ 

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More Notations

First Computations Along Came Poisson Wrap It Up, David

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  $\times \sum_{(\alpha,\beta,\xi,\eta)\in\Omega} \left[ \left(ct_1\alpha^{(1)}\right)^2 + \left(ct_2\alpha^{(2)}\right)^2 + \left(\frac{\xi}{t_1t_2}\right) \right]^{-\frac{3\nu_1}{2}}$   
  $\times \left[ \left(\frac{\beta^{(1)}}{ct_1}\right)^2 + \left(\frac{\beta^{(2)}}{ct_2}\right)^2 + (\eta t_1 t_2)^2 \right]^{-\frac{3\nu_2}{2}}$ 

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More Notations First Computations Along Came Poisso

## The $SL_3(\mathbb{Z})$ Eisenstein series III

### True facts about $G_{\nu_1,\nu_2}$

- $G_{\nu_1,\nu_2}(t_1, t_2)$  is well defined as a converging series when Re  $\nu_1$  and Re  $\nu_2$  are bigger than  $\frac{2}{3}$ ;
- it has meromorphic continuation to  $\mathbb{C}^2$ ;

• 
$$G_{1-\nu_1-\nu_2,\nu_1}(t_1,t_2)=G_{\nu_1,\nu_2}(t_1,t_2).$$

### Proof.

 $G_{
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More Notations First Computations Along Came Poissor Wrap It Up, David

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More Notations First Computations Along Came Poissor Wrap It Up, David

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More Notations First Computations Along Came Poisson Wrap It Up, David

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More Notations First Computations Along Came Poisson Wrap It Up, David

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More Notations First Computations Along Came Poisson Wrap It Up, David

## Action of $\mathfrak{o}^{ imes}$ on $(\mathbb{R}^{\star}_+)^2$

### Definition

If  $\epsilon$  is a unit in  $\mathfrak{o}$  and  $t_1, t_2$  are positive real numbers, we define

 $\epsilon(t_1, t_2) = \left( |\epsilon^{(1)}| t_1, |\epsilon^{(2)}| t_2 \right)$ 

### Remarks

- $|\epsilon^{(1)}\epsilon^{(2)}| = |\mathbb{N}(\epsilon)|$  is a unit in  $\mathbb{Z}$ , so  $|\epsilon^{(2)}| = |\epsilon|^{-1}$ .
- As a consequence,  $\epsilon(t_1, t_2) = (|\epsilon|t_1, |\epsilon|^{-1}t_2)$ .
- By the Unit Theorem, there exists  $\varepsilon \in \mathfrak{o}^{\times}$ , such that

$$\mathfrak{p}^{\times} = \left\{ \pm \varepsilon^n \mid n \in \mathbb{Z} \right\} \quad \text{and} \quad \varepsilon > 1$$

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More Notations First Computations Along Came Poisson Wrap It Up, David

## Action of $\mathfrak{o}^{ imes}$ on $(\mathbb{R}^{\star}_+)^2$

### Definition

If  $\epsilon$  is a unit in  $\mathfrak{o}$  and  $t_1, t_2$  are positive real numbers, we define

$$\epsilon(t_1, t_2) = \left( |\epsilon^{(1)}| t_1, |\epsilon^{(2)}| t_2 \right)$$

### Remarks

- $|\epsilon^{(1)}\epsilon^{(2)}| = |\mathbb{N}(\epsilon)|$  is a unit in  $\mathbb{Z}$ , so  $|\epsilon^{(2)}| = |\epsilon|^{-1}$ .
- As a consequence,  $\epsilon(t_1, t_2) = (|\epsilon|t_1, |\epsilon|^{-1}t_2)$ .
- By the Unit Theorem, there exists  $\varepsilon \in \mathfrak{o}^{\times}$ , such that

$$\mathfrak{o}^{ imes} = ig\{ \pm arepsilon^n \mid n \in \mathbb{Z} ig\} \qquad ext{and} \qquad arepsilon > 1 ig\}$$

More Notations First Computations Along Came Poisson Wrap It Up, David



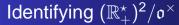
### Theorem

The action of  $\mathfrak{o}^{\times}$  on  $(\mathbb{R}_+^{\star})^2$  has kernel  $\{\pm 1\}.$  A fundamental domain is given by

$$(\mathbb{R}^{\star}_{+})^{2}/\mathfrak{o}^{ imes} = \left\{ (t_{1}, t_{2}) \in (\mathbb{R}^{\star}_{+})^{2} \mid \varepsilon^{-1} \leqslant rac{t_{2}}{t_{1}} \leqslant \varepsilon 
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More Notations First Computations Along Came Poisson Wrap It Up, David



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More Notations First Computations Along Came Poisson Wrap It Up, David

### Action of $\mathfrak{o}^{\times}$ on $\Omega$

Recall that

$$\Omega = \left\{ (\alpha, \beta, \xi, \eta) \in \mathfrak{a} \times (\mathfrak{a}\mathfrak{D})^{-1} \times \mathbb{Z}^2 \mid \begin{cases} (\alpha, \xi) \neq 0\\ (\beta, \eta) \neq 0\\ \operatorname{Tr} \alpha\beta + \xi\eta = 0 \end{cases} \right\}$$

### Definition

If  $\epsilon \in \mathfrak{o}^{\times}$  and  $(\alpha, \beta, \xi, \eta) \in \Omega$ , define

$$\epsilon(\alpha, \beta, \xi, \eta) = (\epsilon \alpha, \epsilon^{-1} \beta, \xi, \eta)$$

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More Notations First Computations Along Came Poisson Wrap It Up, David

## Identifying $\Omega/\mathfrak{o}^{\times}$

Let  $\mathscr{A}$  be a set of representatives of the principal ideals in  $\mathfrak{a}$ . In other words,  $\mathscr{A} = \mathfrak{a}/\mathfrak{o}^{\times}$ .

#### Theorem

The action of  $\mathfrak{o}^{\times}$  on  $\Omega$  is faithful. A complete set of representatives for the orbits is

$$\tilde{\Omega} = \left\{ (\alpha, \beta, \xi, \eta) \in \mathscr{A} \times (\mathfrak{a}\mathfrak{D})^{-1} \times \mathbb{Z} \times \mathbb{Z} \mid \begin{cases} (\alpha, \xi) \neq 0\\ (\beta, \eta) \neq 0\\ Tr \alpha \beta + \xi \eta \end{cases} \right\}$$

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#### More Notations First Computations Along Came Poisson Wrap It Up, David

## The Link

### Define

• 
$$\Lambda = \frac{1}{4} \pi^{-\frac{3\nu_1}{2}} \Gamma\left(\frac{3\nu_1}{2}\right) \pi^{-\frac{3\nu_2}{2}} \Gamma\left(\frac{3\nu_2}{2}\right) \zeta^* (3\nu_1 + 3\nu_2 + 1) \text{ and}$$
  
• 
$$H(\alpha, \beta, \xi, \eta, t_1, t_2) = (t_1 t_2)^s \left[ (ct_1 \alpha^{(1)})^2 + (ct_2 \alpha^{(2)})^2 + (\frac{\xi}{t_1 t_2}) \right]^{-\frac{3\nu_1}{2}} \times \left[ \left(\frac{\beta^{(1)}}{ct_1}\right)^2 + \left(\frac{\beta^{(2)}}{ct_2}\right)^2 + (\eta t_1 t_2)^2 \right]^{-\frac{3\nu_2}{2}}$$

Then

• 
$$G_{\nu_1,\nu_2}(t_1,t_2)(t_1t_2)^s = \Lambda \sum_{(\alpha,\beta,\xi,\eta)\in\Omega} H(\alpha,\beta,\xi,\eta,t_1,t_2)$$

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More Notations First Computations Along Came Poisson Wrap It Up, David

## The Link

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More Notations First Computations Along Came Poisson Wrap It Up, David

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More Notations First Computations Along Came Poisson Wrap It Up, David

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More Notations First Computations Along Came Poisson Wrap It Up, David

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More Notations First Computations Along Came Poisson Wrap It Up, David

## The Link II

We want to work with 
$$I = \iint_{(\mathbb{R}^+_+)/\mathfrak{o}^{\times}} G_{\nu_1,\nu_2}(t_1,t_2)(t_1t_2)^s \frac{\mathrm{d}t_1 \mathrm{d}t_2}{t_1t_2}$$
. Using our

two actions of  $\mathfrak{o}^{\times}$ ,

$$I = \Lambda \sum_{(\alpha,\beta,\xi,\eta)\in\Omega} \iint_{(\mathbb{R}^{+}_{+})^{2}/\mathfrak{o}^{\times}} H(\alpha,\beta,\xi,\eta,t_{1},t_{2}) \frac{dt_{1}dt_{2}}{t_{1}t_{2}}$$
$$= \Lambda \sum_{(\alpha,\beta,\xi,\eta)\in\tilde{\Omega}} \sum_{\epsilon\in\mathfrak{o}^{\times}} \iint_{(\mathbb{R}^{+}_{+})^{2}/\mathfrak{o}^{\times}} H(\epsilon(\alpha,\beta,\xi,\eta),t_{1},t_{2}) \frac{dt_{1}dt_{2}}{t_{1}t_{2}}$$
$$= \Lambda \sum_{(\alpha,\beta,\xi,\eta)\in\tilde{\Omega}} \sum_{\epsilon\in\mathfrak{o}^{\times}} \iint_{\epsilon(\mathbb{R}^{+}_{+})^{2}/\mathfrak{o}^{\times}} H(\alpha,\beta,\xi,\eta,t_{1},t_{2}) \frac{dt_{1}dt_{2}}{t_{1}t_{2}}$$

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## The Link II

More Notations First Computations Along Came Poisson Wrap It Up, David

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## The Link II

More Notations First Computations Along Came Poisson Wrap It Up, David

We want to work with 
$$I = \iint_{(\mathbb{R}_+^*)/\mathfrak{o}^{\times}} G_{\nu_1,\nu_2}(t_1,t_2)(t_1t_2)^s \frac{\mathrm{d}t_1 \mathrm{d}t_2}{t_1t_2}$$
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## The Link II

More Notations First Computations Along Came Poisson Wrap It Up, David

We want to work with 
$$I = \iint_{(\mathbb{R}^*_+)/\mathfrak{o}^{\times}} G_{\nu_1,\nu_2}(t_1,t_2)(t_1t_2)^s \frac{\mathrm{d}t_1 \mathrm{d}t_2}{t_1t_2}$$
. Using our

two actions of  $\mathfrak{o}^{\times}$ ,

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## The Link III

 $I = 2\Lambda \sum_{(\alpha,\beta,\xi,\eta)\in\tilde{\Omega}} \iint_{(\mathbb{R}^{*}_{+})^{2}} H(\alpha,\beta,\xi,\eta,t_{1},t_{2}) \frac{\mathrm{d}t_{1}\mathrm{d}t_{2}}{t_{1}t_{2}}$ 

There are issues with this computation, though:

- We were not careful about justifying interverting integrals and sums.
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$$H(0,0,1,1,t_1,t_2) = (t_1t_2)^{s+3\nu_1-3\nu_2}$$

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# **Troop Reorganization**

We find it necessary to partition  $\Omega$  into smaller subsets and sort out who is integrable and who is not.

The right way to do this is by conditioning the parameters  $\alpha$ ,  $\beta$ ,  $\xi$ ,  $\eta$  on whether they are 0 or not.

		$\alpha = 0  \xi \neq 0$	$\alpha \neq 0  \xi = 0$	$\alpha \neq 0  \xi \neq 0$
$\beta = 0$	$\eta \neq 0$	Excluded	$\Omega_6$	Excluded
eta  eq <b>0</b>	$\eta = 0$	$\Omega_5$	$\Omega_4$	$\Omega_2$
eta  eq <b>0</b>	$\eta \neq 0$	Excluded	$\Omega_3$	$\Omega_1$

Notice that each of these sets is stable under the action of  $\mathfrak{o}^{\times}$ . We call  $G_{\nu_1,\nu_2}^{\Omega_i}$  the part of  $G_{\nu_1,\nu_2}$  corresponding to  $\Omega_i$ .

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# Finally, Some Progress

#### Proposition

Let  $s_0$ ,  $s_1$ ,  $s_2$ ,  $\nu_1$ ,  $\nu_2$  and s be complex numbers related by

$$\begin{cases} 3\nu_1 = s_0 + s_1 + 2s_2 - 1\\ 3\nu_2 = s_0 - s_1 - 2s_2 + 2\\ s = -2s_1 + 2s_2 \end{cases} \iff \begin{cases} \frac{3\nu_1 + 3\nu_2}{2} = s_0 + \frac{1}{2}\\ \frac{3\nu_1 - 3\nu_2 - 2s}{6} = s_1 - \frac{1}{2}\\ \frac{3\nu_1 - 2\nu_2 + s}{6} = s_2 - \frac{1}{2} \end{cases}$$

Assume that s in is some bounded open set, Re  $\nu_1$ , Re  $\nu_2$  and Re  $(\nu_1 - \nu_2)$  are large. Then

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# Finally Some Progress II

### Proposition

$$\begin{split} \iint_{(\mathbb{R}^{*}_{+})^{2}/\mathfrak{o}^{\times}} G^{\Omega_{1}}_{\nu_{1},\nu_{2}}(t_{1},t_{2})(t_{1}t_{2})^{s} \frac{\mathrm{d}t_{1}\mathrm{d}t_{2}}{t_{1}t_{2}} &= \frac{4}{3}\zeta^{\star}(2s_{0}) \\ & \times \sum_{\substack{\xi \in \mathfrak{D}^{-1}, n \in \mathbb{Z} \\ \xi, n \neq 0 \\ \mathrm{Tr}\xi + n = 0}} \tau_{s_{1}-\frac{1}{2}}(|n|)\tau^{K,A}_{s_{2}-\frac{1}{2}}(\xi\mathfrak{D}) \\ & \int_{0}^{+\infty} y^{\frac{3}{2}} \mathcal{K}_{s_{1}-\frac{1}{2}}(2\pi y|n|)\mathcal{K}_{s_{2}-\frac{1}{2}}(2\pi y|\xi^{(1)}|)\mathcal{K}_{s_{2}-\frac{1}{2}}(2\pi y|\xi^{(2)}|)y^{s_{0}-1}\frac{\mathrm{d}y}{y} \end{split}$$

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More Progress

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### Proposition

### Same hypotheses as before. We have

$$\begin{split} &\iint_{\substack{(\mathbb{R}^{*}_{+})^{2}/\mathfrak{o}^{\times}}} G^{\Omega_{2}}_{\nu_{1},\nu_{2}}(t_{1},t_{2})(t_{1}t_{2})^{s} \frac{dt_{1}dt_{2}}{t_{1}t_{2}} &= \frac{2}{3}\zeta^{*}(2s_{0})\zeta^{*}(2-2s_{1})\\ &\times \sum_{\substack{\xi \in \mathfrak{D}^{-1}\\\xi \neq 0}} \tau^{K,A}_{s_{2}-\frac{1}{2}}(\xi\mathfrak{D}) \int_{0}^{+\infty} y^{2-s_{1}} \mathcal{K}_{s_{2}-\frac{1}{2}}(2\pi y|\xi^{(1)}|) \mathcal{K}_{s_{2}-\frac{1}{2}}(2\pi y|\xi^{(2)}|) y^{s_{0}-1} \frac{dy}{y} \end{split}$$

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# The Case of $\Omega_3$

- Unfortunately, we run into trouble when trying to do a similar computation with G<sup>Ω3</sup><sub>ν1,ν2</sub>.
- More precisely, along the way, we encounter

$$\int_{0}^{+\infty} e^{-x} x^{-\frac{3\nu_1 - 3\nu_2 - 2s}{6}} \frac{dx}{x}$$

• This requires Re  $\frac{3\nu_1-3\nu_2-2s}{6}$  to be negative, which contradicts the hypotheses made so far.

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### A Workaround: Poisson's Summation Formula

#### Poisson's Summation Formula

$$\forall t > 0 \qquad \sum_{\eta \in \mathbb{Z}} e^{-\pi \eta^2 t} = \frac{1}{\sqrt{t}} \sum_{\eta \in \mathbb{Z}} e^{-\frac{\pi \eta^2}{t}}$$

This allows us to transform  $G_{\nu_1,\nu_2}^{\Omega_3\cup\Omega_4}$  and obtain a new expression on which the former strategy works.

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More Notations First Computations Along Came Poisson Wrap It Up, David

### A Workaround: Poisson's Summation Formula

### Poisson's Summation Formula

$$orall t > 0 \qquad \sum_{\eta \in \mathbb{Z}} \mathrm{e}^{-\pi \eta^2 t} = rac{1}{\sqrt{t}} \sum_{\eta \in \mathbb{Z}} \mathrm{e}^{-rac{\pi \eta^2}{t}}$$

This allows us to transform  $G_{\nu_1,\nu_2}^{\Omega_3\cup\Omega_4}$  and obtain a new expression on which the former strategy works.

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# A Workaround: Poisson's Summation Formula

#### Poisson's Summation Formula

$$\forall t > 0$$
  $\sum_{\eta \in \mathbb{Z}} \mathbf{e}^{-\pi \eta^2 t} = \frac{1}{\sqrt{t}} \sum_{\eta \in \mathbb{Z}} \mathbf{e}^{-\frac{\pi \eta^2}{t}}$ 

This allows us to transform  $G_{\nu_1,\nu_2}^{\Omega_3\cup\Omega_4}$  and obtain a new expression on which the former strategy works.

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### More Precisely...

### Proposition

### Same hypotheses as before. We have

$$\iint_{\substack{(\mathbb{R}^{*}_{+})^{2}/\mathfrak{o}^{\times}\\ \xi \in \mathfrak{D}^{-1}\\ \xi \neq 0}} \left( G_{\nu_{1},\nu_{2}}^{\Omega_{3}\cup\Omega_{4}}(t_{1},t_{2}) - V_{\nu_{1},\nu_{2}}(t_{1},t_{2}) \right) (t_{1}t_{2})^{s} \frac{\mathrm{d}t_{1}\mathrm{d}t_{2}}{t_{1}t_{2}} = \frac{2}{3}\zeta^{*}(2s_{0})\zeta^{*}(2s_{1})$$

$$\times \sum_{\substack{\xi \in \mathfrak{D}^{-1}\\ \xi \neq 0}} \tau_{s_{2}-\frac{1}{2}}^{K,A}(\xi\mathfrak{D}) \int_{0}^{+\infty} y^{s_{1}+1} K_{s_{2}-\frac{1}{2}}(2\pi y|\xi^{(1)}|) K_{s_{2}-\frac{1}{2}}(2\pi y|\xi^{(2)}|) y^{s_{0}-1} \frac{\mathrm{d}y}{y}$$

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### More Precisely II

### Proposition

### where

$$V_{\nu_1,\nu_2}(t_1,t_2) = \zeta^* (3\nu_1) \zeta^* (3\nu_2 - 1) (t_1 t_2)^{-\frac{3\nu_1 - 3\nu_2 + 3}{2}} \\ E^* \left( \frac{t_1 \alpha_1^{(1)} - i t_2 \alpha_1^{(2)}}{t_1 \alpha_2^{(1)} - i t_2 \alpha_2^{(2)}}, \frac{3\nu_1 + 3\nu_2 - 1}{2} \right)$$

and  $(\alpha_1, \alpha_2)$  is a  $\mathbb{Z}$ -basis of  $\mathfrak{a}$ .

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# What about $\Omega_5$ and $\Omega_6$ ?

It can be shown that

$$\begin{split} G_{\nu_1,\nu_2}^{\Omega_5}(t_1,t_2) &= (t_1 t_2)^{3\nu_1 + \frac{3\nu_2}{2}} \zeta^*(3\nu_1) \zeta^*(3\nu_1 + 3\nu_2 - 1) \\ &\times E^* \bigg( \frac{t_1 \alpha_1^{(1)} + \mathrm{i} t_2 \alpha_1^{(2)}}{t_1 \alpha_2^{(1)} + \mathrm{i} t_2 \alpha_2^{(2)}}, \frac{3\nu_2}{2} \bigg) \end{split}$$

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# In Summary

#### Letting

$$\begin{split} f_{\nu_{1},\nu_{2}}(t_{1},t_{2}) &= \\ \zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{2}-1)(t_{1}t_{2})^{-\frac{3\nu_{1}-3\nu_{2}+3}{2}}E^{*}\bigg(\frac{t_{1}\alpha_{1}^{(1)}-it_{2}\alpha_{1}^{(2)}}{t_{1}\alpha_{2}^{(1)}-it_{2}\alpha_{2}^{(2)}},\frac{3\nu_{1}+3\nu_{2}-1}{2}\bigg) \\ &+(t_{1}t_{2})^{3\nu_{1}+\frac{3\nu_{2}}{2}}\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}+3\nu_{1}-1)E^{*}\bigg(\frac{t_{1}\alpha_{1}^{(1)}+it_{2}\alpha_{1}^{(2)}}{t_{1}\alpha_{2}^{(1)}+it_{2}\alpha_{2}^{(2)}},\frac{3\nu_{2}}{2}\bigg) \\ &+(t_{1}t_{2})^{-3\nu_{2}-\frac{3\nu_{1}}{2}}\zeta^{*}(3\nu_{2})\zeta^{*}(3\nu_{1}+3\nu_{2}-1)E^{*}\bigg(\frac{t_{1}\alpha_{1}^{(1)}+it_{2}\alpha_{1}^{(2)}}{t_{1}\alpha_{2}^{(1)}+it_{2}\alpha_{2}^{(2)}},\frac{3\nu_{1}}{2}\bigg) \end{split}$$

we obtain

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# In Summary II

### Theorem

Same hypotheses as before. We have

$$R_{A}(s_{0}, s_{1}, s_{2}) = 6 \iint_{(\mathbb{R}^{*}_{+})^{2}/\mathfrak{o}^{\times}} (G_{\nu_{1}, \nu_{2}}(t_{1}, t_{2}) - f_{\nu_{1}, \nu_{2}}(t_{1}, t_{2}))(t_{1}t_{2})^{s} \frac{dt_{1}dt_{2}}{t_{1}t_{2}}$$

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### What have we got so far?

- So, we have an identity linking our renormalized integral  $R_A$ ,  $G_{\nu_1,\nu_2}$  and  $f_{\nu_1,\nu_2}$ .
- Although *w* is a functional equation for  $G_{\nu_1,\nu_2}$ , it does not leave  $f_{\nu_1,\nu_2}$  invariant.
- But, as it turns out, the integration over (ℝ<sup>\*</sup><sub>+</sub>)<sup>2</sup>/o<sup>×</sup> has the effect of creating this extra functional equation.
- This requires some work, though.

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More Notations First Computations Along Came Poisson Wrap It Up, David

### What have we got so far?

- So, we have an identity linking our renormalized integral  $R_A$ ,  $G_{\nu_1,\nu_2}$  and  $f_{\nu_1,\nu_2}$ .
- Although *w* is a functional equation for  $G_{\nu_1,\nu_2}$ , it does not leave  $f_{\nu_1,\nu_2}$  invariant.
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### Hmm... Kinda nasty...

$$\frac{R_{A}(s_{0}, s_{1}, s_{2})}{3} = \int_{0}^{+\infty} \left( \mathcal{G}_{\nu_{1}, \nu_{2}}(v) - v^{-\frac{3\nu_{1} - 3\nu_{2} + 3}{2}} L\left(\frac{3\nu_{1} + 3\nu_{2} - 1}{2}\right) \zeta^{\star}(3\nu_{1}) \zeta^{\star}(3\nu_{2} - 1) - v^{3\nu_{1} + \frac{3\nu_{2}}{2}} L\left(\frac{3\nu_{2}}{2}\right) \zeta^{\star}(3\nu_{1}) \zeta^{\star}(3\nu_{1} + 3\nu_{2} - 1) - v^{-3\nu_{2} - \frac{3\nu_{1}}{2}} L\left(\frac{3\nu_{1}}{2}\right) \zeta^{\star}(3\nu_{2}) \zeta^{\star}(3\nu_{1} + 3\nu_{2} - 1) \right) v^{s} \frac{dv}{v}$$

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# **Final Modifications**

• Split this integral as 
$$\int_{0}^{1} + \int_{1}^{+\infty}$$
.

- Replace v by  $\frac{1}{v}$  in the first one.
- Regroup into one integral from 1 to  $+\infty$ .
- Let  $x = \sqrt{v}$ , in order to get rid of those fractional powers.
- Shake.

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## Nastier

$$\begin{aligned} \frac{R_{A}(s_{0},s_{1},s_{2})}{3} &= \int_{1}^{+\infty} \left( \mathcal{G}_{\nu_{1},\nu_{2}}(x^{2})x^{2s} + \mathcal{G}_{\nu_{1},\nu_{2}}(x^{-2})x^{-2s} \\ &- x^{-(3\nu_{1}+3\nu_{2}+3-2s)}L(\frac{3\nu_{1}+3\nu_{2}-1}{2})\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}-1) \\ &- x^{6\nu_{1}+3\nu_{2}+2s}L(\frac{3\nu_{2}}{2})\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}+3\nu_{2}-1) \\ &- x^{-(3\nu+6\nu_{2}-2s)}L(\frac{3\nu_{1}}{2})\zeta^{*}(3\nu_{2})\zeta^{*}(3\nu_{1}+3\nu_{2}-1) \\ &- x^{3\nu_{1}-3\nu_{2}+3-2s}L(\frac{3\nu_{1}+3\nu_{2}-1}{2})\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}-2) \\ &- x^{-(6\nu_{1}+3\nu_{2}+2s)}L(\frac{3\nu_{2}}{2})\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}+3\nu_{2}-1) \\ &- x^{3\nu_{1}+6\nu_{2}-2s}L(\frac{3\nu_{1}}{2})\zeta^{*}(3\nu_{2})\zeta^{*}(3\nu_{1}+3\nu_{2}-1) \right) \frac{dx}{x} \end{aligned}$$

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# **Even Nastier**

$$\begin{split} \frac{R_{A}(s_{0},s_{1},s_{2})}{3} &= \int_{1}^{+\infty} \left( \mathcal{G}_{\nu_{1},\nu_{2}}(x^{2})x^{2s} + \mathcal{G}_{\nu_{1},\nu_{2}}(x^{-2})x^{-2s} \\ &- x^{6\nu_{1}+3\nu_{2}+2s}L(\frac{3\nu_{2}}{2})\zeta^{\star}(3\nu_{1})\zeta^{\star}(3\nu_{1}+3\nu_{2}-1) \\ &- x^{3\nu_{1}-3\nu_{2}+3-2s}L(\frac{3\nu_{1}+3\nu_{2}-1}{2})\zeta^{\star}(3\nu_{1})\zeta^{\star}(3\nu_{1}-2) \\ &- x^{3\nu_{1}+6\nu_{2}-2s}L(\frac{3\nu_{1}}{2})\zeta^{\star}(3\nu_{2})\zeta^{\star}(3\nu_{1}+3\nu_{2}-1) \right) \frac{dx}{x} \\ &- \frac{1}{3\nu_{1}-3\nu_{2}+3-2s}L(\frac{3\nu_{1}+3\nu_{2}-1}{2})\zeta^{\star}(3\nu_{1})\zeta^{\star}(3\nu_{2}-1) \\ &- \frac{1}{3\nu_{1}+6\nu_{2}-2s}L(\frac{3\nu_{1}}{2})\zeta^{\star}(3\nu_{2})\zeta^{\star}(3\nu_{1}+3\nu_{2}-1) \\ &- \frac{1}{6\nu_{1}+3\nu_{2}+2s}L(\frac{3\nu_{2}}{2})\zeta^{\star}(3\nu_{1})\zeta^{\star}(3\nu_{1}+3\nu_{2}-1) \end{split}$$

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$$\begin{split} \frac{R_A(s_0,s_1,s_2)}{3} &= \int\limits_{1}^{+\infty} \left( \mathcal{G}_{\nu_1,\nu_2}(x^2) x^{2s} + \mathcal{G}_{\nu_1,\nu_2}(x^{-2}) x^{-2s} \\ &- x^{6\nu_1 + 3\nu_2 + 2s} L\left(\frac{3\nu_2}{2}\right) \zeta^*(3\nu_1) \zeta^*(3\nu_1 + 3\nu_2 - 1) \\ &- x^{3\nu_1 - 3\nu_2 + 3 - 2s} L\left(\frac{3\nu_1 + 3\nu_2 - 1}{2}\right) \zeta^*(3\nu_1) \zeta^*(3\nu_1 - 2) \\ &- x^{3\nu_1 + 6\nu_2 - 2s} L\left(\frac{3\nu_1}{2}\right) \zeta^*(3\nu_2) \zeta^*(3\nu_1 + 3\nu_2 - 1) \right) \frac{dx}{x} \\ &- \frac{1}{3\nu_1 - 3\nu_2 + 3 - 2s} L\left(\frac{3\nu_1 + 3\nu_2 - 1}{2}\right) \zeta^*(3\nu_1) \zeta^*(3\nu_2 - 1) \\ &- \frac{1}{3\nu_1 + 6\nu_2 - 2s} L\left(\frac{3\nu_1}{2}\right) \zeta^*(3\nu_2) \zeta^*(3\nu_1 + 3\nu_2 - 1) \\ &- \frac{1}{6\nu_1 + 3\nu_2 + 2s} L\left(\frac{3\nu_2}{2}\right) \zeta^*(3\nu_1) \zeta^*(3\nu_1 + 3\nu_2 - 1) \end{split}$$

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## **Even More Nastier**

$$\frac{R_{A}(s_{0},s_{1},s_{2})}{3} = \int_{1}^{+\infty} \left( \mathcal{G}_{\nu_{1},\nu_{2}}(x^{2})x^{2s} + \mathcal{G}_{\nu_{1},\nu_{2}}(x^{-2})x^{-2s} - x^{6\nu_{1}+3\nu_{2}+2s}L(\frac{3\nu_{2}}{2})\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}+3\nu_{2}-1) - x^{3\nu_{1}-3\nu_{2}+3-2s}L(\frac{3\nu_{1}+3\nu_{2}-1}{2})\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}-2) - x^{3\nu_{1}+6\nu_{2}-2s}L(\frac{3\nu_{1}}{2})\zeta^{*}(3\nu_{2})\zeta^{*}(3\nu_{2}-1)\zeta^{*}(3\nu_{1}+3\nu_{2}-1) - x^{-(3\nu_{1}+6\nu_{2}-6-2s)}L(\frac{3\nu_{1}}{2})\zeta^{*}(3\nu_{2}-1)\zeta^{*}(3\nu_{1}+3\nu_{2}-2) - x^{-(3\nu_{1}-3\nu_{2}-3-2s)}L(\frac{3\nu_{1}+3\nu_{2}-1}{2})\zeta^{*}(3\nu_{1}-1)\zeta^{*}(3\nu_{2}) - x^{-(6\nu_{1}+3\nu_{2}-6+2s)}L(\frac{3\nu_{2}}{2})\zeta^{*}(3\nu_{1}+3\nu_{2}-2)\zeta^{*}(3\nu_{1}-1)\right)\frac{dx}{x}$$

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# **Even More Nastier II**

$$\begin{split} &-\frac{1}{3\nu_1 - 3\nu_2 + 3 - 2s}L\left(\frac{3\nu_1 + 3\nu_2 - 1}{2}\right)\zeta^*(3\nu_1)\zeta^*(3\nu_2 - 1) \\ &-\frac{1}{3\nu_1 + 6\nu_2 - 2s}L\left(\frac{3\nu_1}{2}\right)\zeta^*(3\nu_2)\zeta^*(3\nu_1 + 3\nu_2 - 1) \\ &-\frac{1}{6\nu_1 + 3\nu_2 + 2s}L\left(\frac{3\nu_2}{2}\right)\zeta^*(3\nu_1)\zeta^*(3\nu_1 + 3\nu_2 - 1) \\ &+\frac{1}{3\nu_1 + 6\nu_2 - 6 - 2s}L\left(\frac{3\nu_1}{2}\right)\zeta^*(3\nu_2 - 1)\zeta^*(3\nu_1 + 3\nu_2 - 2) \\ &+\frac{1}{3\nu_1 - 3\nu_2 - 3 - 2s}L\left(\frac{3\nu_1 + 3\nu_2 - 1}{2}\right)\zeta^*(3\nu_1 - 1)\zeta^*(3\nu_2) \\ &+\frac{1}{6\nu_1 + 3\nu_2 + 2s - 6}L\left(\frac{3\nu_2}{2}\right)\zeta^*(3\nu_1 + 3\nu_2 - 2)\zeta^*(3\nu_1 - 1). \end{split}$$

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## **Even More Nastier**

$$\frac{R_{A}(s_{0},s_{1},s_{2})}{3} = \int_{1}^{+\infty} \left( \mathcal{G}_{\nu_{1},\nu_{2}}(x^{2})x^{2s} + \mathcal{G}_{\nu_{1},\nu_{2}}(x^{-2})x^{-2s} - x^{6\nu_{1}+3\nu_{2}+2s}L(\frac{3\nu_{2}}{2})\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}+3\nu_{2}-1) - x^{3\nu_{1}-3\nu_{2}+3-2s}L(\frac{3\nu_{1}+3\nu_{2}-1}{2})\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}-2) - x^{3\nu_{1}+6\nu_{2}-2s}L(\frac{3\nu_{1}}{2})\zeta^{*}(3\nu_{2})\zeta^{*}(3\nu_{1}+3\nu_{2}-1) - x^{-(3\nu_{1}+6\nu_{2}-6-2s)}L(\frac{3\nu_{1}}{2})\zeta^{*}(3\nu_{2}-1)\zeta^{*}(3\nu_{1}+3\nu_{2}-2) - x^{-(3\nu_{1}-3\nu_{2}-3-2s)}L(\frac{3\nu_{1}+3\nu_{2}-1}{2})\zeta^{*}(3\nu_{1}-1)\zeta^{*}(3\nu_{2}) - x^{-(6\nu_{1}+3\nu_{2}-6+2s)}L(\frac{3\nu_{2}}{2})\zeta^{*}(3\nu_{1}+3\nu_{2}-2)\zeta^{*}(3\nu_{1}-1)\right)\frac{dx}{x}$$

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More Notations First Computations Along Came Poisson Wrap It Up, David

## **Even More Nastier**

$$\frac{R_{A}(s_{0},s_{1},s_{2})}{3} = \int_{-\infty}^{+\infty} \left( \mathcal{G}_{\nu_{1},\nu_{2}}(x^{2})x^{2s} + \mathcal{G}_{\nu_{1},\nu_{2}}(x^{-2})x^{-2s} - x^{6\nu_{1}+3\nu_{2}+2s}L(\frac{3\nu_{2}}{2})\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}+3\nu_{2}-1) + x^{3\nu_{1}-3\nu_{2}+3-2s}L(\frac{3\nu_{1}+3\nu_{2}-1}{2})\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}-2) - x^{3\nu_{1}+6\nu_{2}-2s}L(\frac{3\nu_{1}}{2})\zeta^{*}(3\nu_{2})\zeta^{*}(3\nu_{2}+3\nu_{2}-1) - x^{-(3\nu_{1}+6\nu_{2}-6-2s)}L(\frac{3\nu_{1}}{2})\zeta^{*}(3\nu_{2}-1)\zeta^{*}(3\nu_{1}+3\nu_{2}-2) - x^{-(3\nu_{1}-3\nu_{2}-3-2s)}L(\frac{3\nu_{1}+3\nu_{2}-1}{2})\zeta^{*}(3\nu_{1}-1)\zeta^{*}(3\nu_{2}) - x^{-(6\nu_{1}+3\nu_{2}-6+2s)}L(\frac{3\nu_{2}}{2})\zeta^{*}(3\nu_{1}+3\nu_{2}-2)\zeta^{*}(3\nu_{1}-1)\right)\frac{dx}{x}$$

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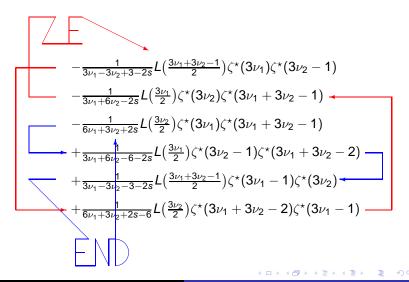
## **Even More Nastier**

$$\frac{R_{A}(s_{0},s_{1},s_{2})}{3} = \int_{-\infty}^{+\infty} \left( \mathcal{G}_{\nu_{1},\nu_{2}}(x^{2})x^{2s} + \mathcal{G}_{\nu_{1},\nu_{2}}(x^{-2})x^{-2s} - x^{6\nu_{1}+3\nu_{2}+2s}L(\frac{3\nu_{2}}{2})\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}+3\nu_{2}-1) - x^{3\nu_{1}-3\nu_{2}+3-2s}L(\frac{3\nu_{1}+3\nu_{2}-1}{2})\zeta^{*}(3\nu_{1})\zeta^{*}(3\nu_{1}-2) - x^{3\nu_{1}+6\nu_{2}-2s}L(\frac{3\nu_{1}}{2})\zeta^{*}(3\nu_{2})\zeta^{*}(3\nu_{1}+3\nu_{2}-1) - x^{-(3\nu_{1}+6\nu_{2}-6-2s)}L(\frac{3\nu_{1}}{2})\zeta^{*}(3\nu_{2}-1)\zeta^{*}(3\nu_{1}+3\nu_{2}-2) - x^{-(3\nu_{1}-3\nu_{2}-3-2s)}L(\frac{3\nu_{1}+3\nu_{2}-1}{2})\zeta^{*}(3\nu_{1}-1)\zeta^{*}(3\nu_{2}) - x^{-(6\nu_{1}+3\nu_{2}-6+2s)}L(\frac{3\nu_{2}}{2})\zeta^{*}(3\nu_{1}+3\nu_{2}-2)\zeta^{*}(3\nu_{1}-1)\right)\frac{dx}{x}$$

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## Even More Nastier II



#### Acknowledgements

#### Thanks, yall!

David LECOMTE Hidden Functional Equations

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#### Acknowledgements

#### Thanks, yall!

David LECOMTE Hidden Functional Equations

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