# Hidden Functional Equations for Rankin-Selberg transforms: New Results 

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## Outline

(9) Introduction and First Notations

- The $S L_{2}(\mathbb{Z})$ Eisenstein Series
- The Hilbert Modular Eisenstein Series
- Zagier's Renormalization
- Previous Results and Present Goal
(2) Polar divisor considerations
- Identifying the Polar Divisor of $R_{A}$
- Finding the Group of $\Pi$
- An Upper Bound for $|S|$
(3) Proof of the Main Theorem
- More Notations
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- Along Came Poisson
- Wrap It Up, David
(4) Acknowledgements

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## The $S_{2}(\mathbb{Z})$ Eisenstein Series

- $E^{\star}(z, s)=\frac{1}{2} \pi^{-s} \Gamma(s) \sum_{(m, n) \in \mathbb{Z}^{2}} \frac{y^{s}}{|m z+n|^{2 s}} \quad z \in \mathscr{H} \quad$ Res $>1$
- $E^{+}(z, s)$ is $S L_{2}(\mathbb{Z})$-automorphic in the variable $z$.
- $E^{\star}(z, s)$ is analytic in the variable $s$.


## The $S L_{2}(\mathbb{Z})$ Eisenstein Series

- $E^{\star}(z, s)=\frac{1}{2} \pi^{-s} \Gamma(s) \sum_{\substack{(m, n) \in \mathbb{Z}^{2} \\(m, n) \neq 0}} \frac{y^{s}}{|m z+n|^{2 s}} \quad z \in \mathscr{H} \quad$ Res $>1$
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## The $L_{2}(\mathbb{Z})$ Eisenstein Series

- $E^{\star}(z, s)=\frac{1}{2} \pi^{-s} \Gamma(s) \sum_{\substack{(m, n) \in \mathbb{Z}^{2} \\(m, n) \neq 0}} \frac{y^{s}}{|m z+n|^{2 s}} \quad z \in \mathscr{H} \quad$ Res $>1$
- $E^{\star}(z, s)$ is $S L_{2}(\mathbb{Z})$-automorphic in the variable $z$.
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## Fourier Expansion


where

- $\tau_{\omega}(n)=n^{-\omega} \sum_{\substack{d \mid n \\ d>0}} d^{2 \omega}$ is a divisor sum;
- $K_{\omega}(y)=\frac{1}{2} \int_{0}^{+\infty} \mathrm{e}^{-\frac{y}{2}\left(t+\frac{1}{t}\right)} t^{\omega} \frac{\mathrm{d} t}{t}$ is a Bessel function.


## Fourier Expansion

$$
E^{\star}(z, s)=y^{s} \zeta^{\star}(2 s)+y^{1-s} \zeta^{\star}(2-2 s)+2 \sqrt{y} \sum_{n \neq 0} \tau_{s-\frac{1}{2}}(|n|) K_{s-\frac{1}{2}}(2 \pi|n| y) \mathrm{e}^{2 \pi i n x}
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## What You Need to Know About $\tau_{\omega}$ and $K_{\omega}$

- $\tau_{\omega}(n)=\tau_{-\omega}(n)$ and $K_{\omega}(y)=K_{-\omega}(y)$;
- $K_{\omega}$ has rapid decay:

$$
\forall y>2 \quad\left|K_{\omega}(y)\right| \leqslant M e^{-\frac{y}{2}}
$$

As a consequence, $E^{\star}(z, s)$ is the uniform limit of its Fourier series, can be analytically continued to $\mathbb{C} \backslash\{0,1\}$ and satisfies the functional equation

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\forall z \in \mathscr{H} \quad \forall s \neq 0,1 \quad E^{\star}(z, s)=E^{\star}(z, 1-s)
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## The Setting

- $\mathbb{Q} \hookrightarrow K$ is a totally real field extension with ring of integers $\mathfrak{o}$ and discriminant $D$.
- There are exactly $N$ distinct embeddings $K \hookrightarrow \mathbb{C}$, denoted as $\sigma_{1}, \ldots, \sigma_{N}$.
- For convenience, we write also

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- $A$ is an ideal class in $K$ and $\mathfrak{a}$ an ideal in $A^{-1}$.


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## The Hilbert Modular Eisenstein Series

- $E_{K, A}^{\star}(z, s)=\mathbb{N}(\mathfrak{a})^{2 s} \pi^{-N s} \Gamma(s)^{N} D^{s}$

where $z \in \mathscr{H}$ and Res>1.
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& +2^{N} y^{\frac{N}{2}} \sum_{\substack{\xi \in \mathfrak{D}^{-1} \\
\xi \neq 0}} \tau_{s-\frac{1}{2}}^{K, A}(\xi \mathfrak{D})\left(\prod_{i=1}^{N} K_{s-\frac{1}{2}}\left(2 \pi y\left|\xi^{(i)}\right|\right)\right) \mathrm{e}^{2 \pi i x \operatorname{Tr} \xi}
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## Zagier's Theorem

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Let $F$ be an $S L_{2}(\mathbb{Z})$-automorphic function on $\mathscr{H}$. Suppose that there exists a function $\varphi$ of the form

such that $F(z)-\varphi(y) \underset{y \rightarrow+\infty}{=} o\left(y^{-n}\right)$ for every positive integer $n$. Then


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\varphi(y)=\sum_{i=1}^{\ell} c_{i} y^{\alpha_{i}} \ln ^{n_{i}} y \quad \alpha_{i} \in \mathbb{C} \quad n_{i} \in \mathbb{N}
$$

such that $F(z)-\varphi(y) \underset{y \rightarrow+\infty}{=} o\left(y^{-n}\right)$ for every positive integer $n$. Then

$$
R(s)=R N \int_{S L_{2}(\mathbb{Z}) \backslash \mathscr{H}} E^{\star}(z, s) F(z) \frac{d z}{y^{2}} \overline{d e f} \int_{0}^{+\infty}\left(a_{0}(y)-\varphi(y)\right) y^{s-1} \frac{d y}{y}
$$

## Zagier's Theorem: Ze Continuation

## Theorem

where $a_{0}(y)=\int_{0}^{1} F(x+i y) d x$, is well defined as an absolutely
convergent integral for Res big enough. It has analytic continuation to $\mathbb{C}$, except for poles at 0,1 , the $\alpha_{i}$ 's and the $\left(1-\alpha_{i}\right)$ 's.
Furthermore, $R(s)=R(1-s)$.

## Previous Results

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- $R N \int_{S L_{2}(\mathbb{Z}) \backslash \mathscr{H}} E^{\star}\left(z, s_{0}\right) E_{K, A}^{\star}\left(z, s_{1}\right) \frac{\mathrm{d} z}{y^{2}}$ when $K$ is a totally real cubic field and $A$ is an ideal class;



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field and $A$ is an ideal class;
- $R N \int_{S_{2}(\mathbb{Z}) \backslash \mathscr{H}} E^{\star}\left(z, s_{0}\right) E^{\star}\left(z, s_{1}\right) E^{\star}\left(z, s_{2}\right) E^{\star}\left(z, s_{3}\right) \frac{\mathrm{d} z}{y^{2}}$.

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## Our Goal

## We suspect that a similar phenomenon occurs and want to identify the full group of functional equations for the renormalized Rankin-Selberg transform

$$
R_{A}\left(s_{0}, s_{1}, s_{2}\right)=\int_{s L_{2}(\mathbb{Z}) \backslash \mathscr{H}} E^{\star}\left(z, s_{0}\right) E^{\star}\left(z, s_{1}\right) E_{K, A}^{\star}\left(z, s_{2}\right) \frac{\mathrm{d} z}{y^{2}}
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when $K$ is a real quadratic field and $A$ is an ideal class. Zagier's theorem, together with our knowledge of the objects involved, predict 16 functional equations. In fact,

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## Main Theorem

## Theorem

Let $K$ be a real quadratic field and $A$ an ideal class. The function $R_{A}\left(s_{0}, s_{1}, s_{2}\right)$ has a group of functional equations of order 48. It is generated by the transformations

$$
s_{0} \longmapsto 1-s_{0} \quad s_{1} \longmapsto 1-s_{1} \quad s_{2} \longmapsto 1-s_{2} \quad s_{0} \longleftrightarrow s_{1}
$$

and

$$
\left[\begin{array}{l}
s_{0} \\
s_{1} \\
s_{2}
\end{array}\right] \stackrel{w}{\longmapsto}\left[\begin{array}{l}
-\frac{s_{0}}{2}+\frac{s_{1}}{2}+s_{2} \\
1-\frac{s_{0}}{2}+\frac{s_{1}}{2}-s_{2} \\
1-\frac{s_{0}}{2}-\frac{s_{1}}{2}
\end{array}\right] .
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## Strategy

> The polar divisor $\Pi$ of $R_{A}$ is the subset of $\mathbb{C}^{3}$ at which $R_{A}$ is undefined. Zagier's theorem provides us with a complete description of $\Pi$. All we have to do is identify the part $\varphi$ of $F$ that is not of rapid decay and hope it looks like


Once we know the $\alpha_{i}$ 's, we know $\Pi$.

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## Reminder

## Remember that

$E^{\star}\left(z, s_{1}\right)=y^{s_{1}} \zeta^{\star}\left(2 s_{1}\right)+y^{1-s_{1}} \zeta^{\star}\left(2-2 s_{1}\right)$ $+2 \sqrt{y} \sum_{n \neq 0} \tau_{s_{1}-\frac{1}{2}}(|n|) K_{s_{1}-\frac{1}{2}}(2 \pi|n| y) \mathrm{e}^{2 \pi i n x}$

$$
E_{K, A}^{\star}\left(z, s_{2}\right)=y^{2 s_{2}} \zeta_{K, A}^{\star}\left(2 s_{2}\right)+y^{2\left(1-s_{2}\right)} \zeta_{K, A}^{\star}\left(2-2 s_{2}\right)
$$

$$
+4 y \sum_{\substack{\xi \in \mathfrak{D}-1 \\ \epsilon \pm 0}} \tau_{s_{2}-\frac{1}{2}}^{K, A}(\xi \mathfrak{D})\left(\prod_{i=1}^{2} K_{S_{2}-\frac{1}{2}}\left(2 \pi y\left|\xi^{(i)}\right|\right)\right) \mathrm{e}^{2 \pi i x \operatorname{Tr} \xi}
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## Expression of $\varphi$



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\begin{aligned}
& \varphi(y)=y^{s_{1}+2 s_{2}} \zeta^{\star}\left(2 s_{1}\right) \zeta_{K, A}^{\star}\left(2 s_{2}\right)+y^{2+s_{1}-2 s_{2}} \zeta^{\star}\left(2 s_{1}\right) \zeta_{K, A}^{\star}\left(2-2 s_{2}\right) \\
& +y^{1-s_{1}+2 s_{2}} \zeta^{\star}\left(2-2 s_{1}\right) \zeta_{\kappa, A}^{\star}\left(2 s_{2}\right)+y^{3-s_{1}-2 s_{2}} \zeta^{\star}\left(2-2 s_{1}\right) \zeta_{K, A}^{\star}\left(2-2 s_{2}\right)
\end{aligned}
$$

## Cartesian Equation of $\Pi$

## By Zagier's Theorem, the polar divisor of $R_{A}$ is the union of the 14 hyperplanes

$$
s_{0}=s_{1}+2 s_{2} \quad s_{0}=2+s_{1}-2 s_{2} \quad s_{0}=1-s_{1}+2 s_{2} \quad s_{0}=3-s_{1}-2 s_{2}
$$

$$
s_{0}=1-s_{1}-2 s_{2} \quad s_{0}=-1-s_{1}+2 s_{2} \quad s_{0}=s_{1}-2 s_{2} \quad s_{0}=-2+s_{1}+2 s_{2}
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s_{0}, s_{1}, s_{2}=0 \quad \text { and } \quad s_{0}, s_{1}, s_{2}=1
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## Cartesian Equation of $\Pi$

By Zagier's Theorem, the polar divisor of $R_{A}$ is the union of the 14 hyperplanes
$s_{0}=1-s_{1}-2 s_{2} \quad s_{0}=-1-s_{1}+2 s_{2} \quad s_{0}=s_{1}-2 s_{2} \quad s_{0}=-2+s_{1}+2 s_{2}$

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## So what does П look like?

The 14 hyperplanes in $\mathbb{C}^{3}$ cut out a rhombic dodecahedron:


## The Group of $\Pi$

The group $S_{\Pi}$ of transformations leaving $\Pi$ invariant is computed through standard counting arguments.
For example, $S_{\Pi}$ acts transitively on the set of vertices that are connected to 4 other vertices.
There are 6 such vertices.
Each of them is fixed by 8 transformations:
4 reflexions and 4 rotations.


So $\left|S_{\Pi}\right|=6 \times 8=48$.

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## Relationship Between $S$ and $S_{\Pi}$

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## The $S L_{3}(\mathbb{Z})$ Eisenstein series I

```
- \(\nu_{1}\) and \(\nu_{2}\) are complex numbers with real part bigger than \(\frac{2}{3}\).
- \(\Omega\) is the set of quadruples \((\alpha, \beta, \xi, \eta) \in \mathfrak{a} \times(\mathfrak{a} \mathfrak{D})^{-1} \times \mathbb{Z} \times \mathbb{Z}\) with
- \(\alpha\) and \(\xi\) not simultaneously 0;
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- \(\operatorname{Tr} \alpha \beta+\xi \eta=0\).
- \(c=\mathbb{N}(a)^{-\frac{1}{2}} D^{-\frac{1}{4}}\);
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## The $S L_{3}(\mathbb{Z})$ Eisenstein Series II



## The $S L_{3}(\mathbb{Z})$ Eisenstein Series II

$$
\text { Let } \begin{aligned}
& G_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)= \frac{1}{4} \pi^{-\frac{3 \nu_{1}}{2}} \Gamma\left(\frac{3 \nu_{1}}{2}\right) \pi^{-\frac{3 \nu_{2}}{2}} \Gamma\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}+1\right) \\
& \times \sum_{(\alpha, \beta, \xi, \eta) \in \Omega}\left[\left(c t_{1} \alpha^{(1)}\right)^{2}+\left(c t_{2} \alpha^{(2)}\right)^{2}+\left(\frac{\xi}{t_{1} t_{2}}\right)\right]^{-\frac{3 \nu_{1}}{2}} \\
& \times\left[\left(\frac{\beta^{(1)}}{c t_{1}}\right)^{2}+\left(\frac{\beta^{(2)}}{c t_{2}}\right)^{2}+\left(\eta t_{1} t_{2}\right)^{2}\right]^{-\frac{3 \nu_{2}}{2}}
\end{aligned}
$$

## The $S_{3}(\mathbb{Z})$ Eisenstein series III

## True facts about $G_{\nu_{1}, \nu_{2}}$

- $G_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)$ is well defined as a converging series when $\operatorname{Re} \nu_{1}$ and $\operatorname{Re} \nu_{2}$ are bigger than $\frac{2}{3}$;
- it has meromorphic continuation to $\mathbb{C}^{2}$;
- $G_{1-\nu_{1}-\nu_{2}, \nu_{1}}\left(t_{1}, t_{2}\right)=G_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)$.


## Proof.

$G_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)$ is a particular value of the $S L_{3}(\mathbb{Z})$ Eisenstein series.

## The $S L_{3}(\mathbb{Z})$ Eisenstein series III

## True facts about $G_{\nu_{1}, \nu_{2}}$

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## Proof.

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## The $S L_{3}(\mathbb{Z})$ Eisenstein series III

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## The $S_{3}(\mathbb{Z})$ Eisenstein series III

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## Action of $\mathfrak{o}^{\times}$on $\left(\mathbb{R}_{+}^{\star}\right)^{2}$

## Definition

If $\epsilon$ is a unit in 0 and $t_{1}, t_{2}$ are positive real numbers, we define

$$
\epsilon\left(t_{1}, t_{2}\right)=\left(\left|\epsilon^{(1)}\right| t_{1},\left|\epsilon^{(2)}\right| t_{2}\right)
$$

## Remarks

- $\left|\epsilon^{(1)} f^{(2)}\right|=|\mathbb{N}(\epsilon)|$ is a unit in $\mathbb{Z}$, so $\left|\epsilon^{(2)}\right|=|\epsilon|^{-1}$.
- As a consequence, $\epsilon\left(t_{1}, t_{2}\right)=\left(|\epsilon| t_{1},|\epsilon|^{-1} t_{2}\right)$.
- By the Unit Theorem, there exists $\varepsilon \in \mathfrak{o}^{\times}$, such that

$$
0^{x}=\left\{ \pm \varepsilon^{n} \mid n \in \mathbb{Z}\right\} \quad \text { and } \quad \varepsilon>1
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## Identifying $\left(\mathbb{R}_{+}^{\star}\right)^{2} / \mathfrak{o}^{x}$

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The action of ox on $\left(\mathbb{R}^{*}\right)^{2}$ has kernel $\{ \pm 1\}$. A fundamental domain is given by


## Identifying $\left(\mathbb{R}_{+}^{\star}\right)^{2} / \mathfrak{o}^{x}$

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$$
\left(\mathbb{R}_{+}^{\star}\right)^{2} / \mathfrak{o}^{\times}=\left\{\left(t_{1}, t_{2}\right) \in\left(\mathbb{R}_{+}^{\star}\right)^{2} \left\lvert\, \varepsilon^{-1} \leqslant \frac{t_{2}}{t_{1}} \leqslant \varepsilon\right.\right\}
$$

## Action of $\mathfrak{o}^{\times}$on $\Omega$

## Recall that

$$
\Omega=\left\{(\alpha, \beta, \xi, \eta) \in \mathfrak{a} \times(\mathfrak{a} \mathfrak{D})^{-1} \times \mathbb{Z}^{2} \left\lvert\,\left\{\begin{array}{l}
(\alpha, \xi) \neq 0 \\
(\beta, \eta) \neq 0 \\
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\end{array}\right\}\right.\right.
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## Definition

If $\epsilon \in \mathfrak{o}^{\times}$and $(\alpha, \beta, \xi, \eta) \in \Omega$, define
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Introduction and First Notations
Polar divisor considerations
Proof of the Main Theorem Acknowledgements

## Identifying $\Omega / o^{\times}$

Let $\mathscr{A}$ be a set of representatives of the principal ideals in $\mathfrak{a}$. In other words, $\mathscr{A}=\mathfrak{a} / \mathfrak{o}^{\times}$

## Theorem

The action of ${ }^{\times} \times$on $\Omega$ is faithful. A complete set of representatives for the orbits is


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\tilde{\Omega}=\left\{(\alpha, \beta, \xi, \eta) \in \mathscr{A} \times(\mathfrak{a} \mathfrak{D})^{-1} \times \mathbb{Z} \times \mathbb{Z} \left\lvert\,\left\{\begin{array}{l}
(\alpha, \xi) \neq 0 \\
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\end{array}\right\}\right.\right.
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## The Link

Define

- $\wedge=\frac{1}{4} \pi^{-\frac{3 \nu_{1}}{2}} \Gamma\left(\frac{3 \nu_{1}}{2}\right) \pi^{-\frac{3 \nu_{2}}{2}} \Gamma\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}+1\right)$ and
- $H\left(\alpha, \beta, \xi, \eta, t_{1}, t_{2}\right)=\left(t_{1} t_{2}\right)^{s}\left[\left(c t_{1} \alpha^{(1)}\right)^{2}+\left(c t_{2} \alpha^{(2)}\right)^{2}+\left(\frac{\xi}{t_{1} t_{2}}\right)\right]^{-\frac{3 \nu_{1}}{2}}$

Then

$$
\text { - } G_{\cdots, \nu_{2}}\left(t_{1}, t_{2}\right)\left(t_{1} t_{2}\right)^{s}=\Lambda \sum_{(\alpha, \beta, \xi, \eta) \in \Omega} H\left(\alpha, \beta, \xi, \eta, t_{1}, t_{2}\right)
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$$
\times\left[\left(\frac{\beta^{(1)}}{c t_{1}}\right)^{2}+\left(\frac{\beta^{(2)}}{c t_{2}}\right)^{2}+\left(\eta t_{1} t_{2}\right)^{2}\right]^{-\frac{3 \nu_{2}}{2}}
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## The Link

Define

- $\Lambda=\frac{1}{4} \pi^{-\frac{3 \nu_{1}}{2}} \Gamma\left(\frac{3 \nu_{1}}{2}\right) \pi^{-\frac{3 \nu_{2}}{2}} \Gamma\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}+1\right)$ and
- $H\left(\alpha, \beta, \xi, \eta, t_{1}, t_{2}\right)=\left(t_{1} t_{2}\right)^{s}\left[\left(c t_{1} \alpha^{(1)}\right)^{2}+\left(c t_{2} \alpha^{(2)}\right)^{2}+\left(\frac{\xi}{t_{1} t_{2}}\right)\right]^{-\frac{3 \nu_{1}}{2}}$

$$
\times\left[\left(\frac{\beta^{(1)}}{c t_{1}}\right)^{2}+\left(\frac{\beta^{(2)}}{c t_{2}}\right)^{2}+\left(\eta t_{1} t_{2}\right)^{2}\right]^{-\frac{3 \nu_{2}}{2}}
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$$

Then

- $G_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)\left(t_{1} t_{2}\right)^{s}=\Lambda \sum_{(\alpha, \beta, \xi, \eta) \in \Omega} H\left(\alpha, \beta, \xi, \eta, t_{1}, t_{2}\right)$


## The Link II

We want to work with $I=\iint_{\left(\mathbb{R}^{*}\right) / 0^{\times}} G_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)\left(t_{1} t_{2}\right)^{s} \frac{\mathrm{~d} t_{1} \mathrm{~d} t_{2}}{t_{1} t_{2}}$. Using our
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& =\Lambda \sum_{(\alpha, \beta, \xi, \eta) \in \tilde{\Omega} \in \in \mathfrak{o}^{\times}} \sum_{\left(\mathbb{R}_{+}^{*}\right)^{2} / \mathfrak{o}^{\times}} H \int_{\left(\epsilon(\alpha, \beta, \xi, \eta), t_{1}, t_{2}\right) \frac{\mathrm{d} t_{1} \mathrm{~d} t_{2}}{t_{1} t_{2}}}
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\end{aligned}
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## The Link III

$$
I=2 \Lambda \sum_{(\alpha, \beta, \xi, \eta) \in \tilde{\Omega}} \iint_{\left(\mathbb{R}_{+}^{\star}\right)^{2}} H\left(\alpha, \beta, \xi, \eta, t_{1}, t_{2}\right) \frac{\mathrm{d} t_{1} \mathrm{~d} t_{2}}{t_{1} t_{2}}
$$

## There are issues with this computation, though:

- We were not careful about justifying interverting integrals and sums.
- In fact, it is not valid at all: some terms in the sum are not even integrable on $\left(\mathbb{R}_{+}^{\star}\right)^{2} / \mathfrak{o}^{\times}$. For example,
$H\left(0,0,1,1, t_{1}, t_{2}\right)=\left(t_{1} t_{2}\right)^{s+3 \nu_{1}-3 \nu_{2}}$
cannot be integrable at the same time at 0 and $\infty$.


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## Troop Reorganization

We find it necessary to partition $\Omega$ into smaller subsets and sort out who is integrable and who is not.
The right way to do this is by conditioning the parameters $\alpha, \beta, \xi, \eta$ on whether they are 0 or not.

|  | $\alpha=0 \quad \xi \neq 0$ | $\alpha \neq 0 \quad \xi=0$ | $\alpha \neq 0 \quad \xi \neq 0$ |  |
| :---: | :---: | :---: | :---: | :---: |
| $\beta=0$ | $\eta \neq 0$ | Excluded | $\Omega_{6}$ | Excluded |
| $\beta \neq 0$ | $\eta=0$ | $\Omega_{5}$ | $\Omega_{4}$ | $\Omega_{2}$ |
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Notice that each of these sets is stable under the action of o We call $G_{\nu_{1}, \nu_{2}}^{\Omega_{i}}$ the part of $G_{\nu_{1}, \nu_{2}}$ corresponding to $\Omega_{i}$.

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## Finally, Some Progress

## Proposition

Let $s_{0}, s_{1}, s_{2}, \nu_{1}, \nu_{2}$ and $s$ be complex numbers related by


Assume that $s$ in is some bounded open set, $\operatorname{Re} \nu_{1}, \operatorname{Re} \nu_{2}$ and $\operatorname{Re}\left(\nu_{1}-\nu_{2}\right)$ are large. Then

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$$
\left\{\begin{array} { r l } 
{ 3 \nu _ { 1 } = s _ { 0 } + s _ { 1 } + 2 s _ { 2 } - 1 } \\
{ 3 \nu _ { 2 } = s _ { 0 } - s _ { 1 } - 2 s _ { 2 } + 2 } \\
{ s } & { = - 2 s _ { 1 } + 2 s _ { 2 } }
\end{array} \Longleftrightarrow \left\{\begin{array}{ll}
\frac{3 \nu_{1}+3 \nu_{2}}{2} & =s_{0}+\frac{1}{2} \\
\frac{3 \nu_{1}-3 \nu_{2}-2 s}{6} & =s_{1}-\frac{1}{2} \\
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## Finally Some Progress II

## Proposition

$$
\begin{aligned}
& \iint_{\left(\mathbb{R}_{+}^{\star}\right)^{2} / \mathfrak{o}^{\times}} G_{\nu_{1}, \nu_{2}}^{\Omega_{1}}\left(t_{1}, t_{2}\right)\left(t_{1} t_{2}\right)^{s} \frac{\mathrm{~d} t_{1} \mathrm{~d} t_{2}}{t_{1} t_{2}}= \frac{4}{3} \zeta^{\star}\left(2 s_{0}\right) \\
& \times \sum_{\substack{\xi \in \mathfrak{D}^{-1}, n \in \mathbb{Z} \\
\xi, n \neq 0 \\
T r \xi+n=0}} \tau_{s_{1}-\frac{1}{2}}(|n|) \tau_{s_{2}-\frac{1}{2}}^{K, A}(\xi \mathfrak{D}) \\
& \quad \int_{0}^{+\infty} y^{\frac{3}{2}} K_{s_{1}-\frac{1}{2}}(2 \pi y|n|) K_{s_{2}-\frac{1}{2}}\left(2 \pi y\left|\xi^{(1)}\right|\right) K_{s_{2}-\frac{1}{2}}\left(2 \pi y\left|\xi^{(2)}\right|\right) y^{s_{0}-1} \frac{\mathrm{~d} y}{y}
\end{aligned}
$$

## More Progress

## Proposition

Same hypotheses as before. We have

$$
\begin{aligned}
& \quad \iint_{\left(\mathbb{R}_{+}^{\star}\right)^{2} / \mathfrak{o}^{\times}} G_{\nu_{1}, \nu_{2}}^{\Omega_{2}}\left(t_{1}, t_{2}\right)\left(t_{1} t_{2}\right)^{s} \frac{\mathrm{~d} t_{1} \mathrm{~d} t_{2}}{t_{1} t_{2}}=\frac{2}{3} \zeta^{\star}\left(2 s_{0}\right) \zeta^{\star}\left(2-2 s_{1}\right) \\
& \times \sum_{\substack{\xi \in \mathfrak{D}^{-1} \\
\xi \neq 0}} \tau_{s_{2}-\frac{1}{2}}^{K, A}(\xi \mathfrak{D}) \int_{0}^{+\infty} y^{2-s_{1}} K_{s_{2}-\frac{1}{2}}\left(2 \pi y\left|\xi^{(1)}\right|\right) K_{s_{2}-\frac{1}{2}}\left(2 \pi y\left|\xi^{(2)}\right|\right) y^{s_{0}-1} \frac{\mathrm{~d} y}{y}
\end{aligned}
$$

## The Case of $\Omega_{3}$

- Unfortunately, we run into trouble when trying to do a similar computation with $G_{\nu_{1}, \nu_{2}}^{\Omega_{3}}$.
- More precisely, along the way, we encounter

- This requires $\operatorname{Re} \frac{3 \nu_{1}-3 \nu_{2}-2 s}{6}$ to be negative, which contradicts the hypotheses made so far.


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## A Workaround: Poisson's Summation Formula

## Poisson's Summation Formula

This allows us to transform $G_{\nu_{1}, \nu_{2}}^{\Omega_{3} \cup \Omega_{4}}$ and obtain a new expression on which the former strategy works.

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$$
\forall t>0 \quad \sum_{\eta \in \mathbb{Z}} \mathrm{e}^{-\pi \eta^{2} t}=\frac{1}{\sqrt{t}} \sum_{\eta \in \mathbb{Z}} \mathrm{e}^{-\frac{\pi \eta^{2}}{t}}
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This allows us to transform $G_{\nu_{1}, \nu_{2}}^{\Omega_{3} \cup \Omega_{4}}$ and obtain a new expression on which the former strategy works.

## More Precisely...

## Proposition

## Same hypotheses as before. We have

$$
\begin{aligned}
& \iint_{\left(\mathbb{R}_{+}^{\star}\right)^{2} / 0^{\times}}\left(G_{\nu_{1}, \nu_{2}}^{\Omega_{3} \cup \Omega_{4}}\left(t_{1}, t_{2}\right)-V_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)\right)\left(t_{1} t_{2}\right)^{s} \frac{\mathrm{~d} t_{1} \mathrm{~d} t_{2}}{t_{1} t_{2}}=\frac{2}{3} \zeta^{\star}\left(2 s_{0}\right) \zeta^{\star}\left(2 s_{1}\right) \\
& \times \sum_{\substack{\xi \in \mathcal{D}^{-1} \\
\xi \neq 0}} \tau_{s_{2}-\frac{1}{2}}^{K, A}(\xi \mathfrak{D}) \int_{0}^{+\infty} y^{s_{1}+1} K_{s_{2}-\frac{1}{2}}\left(2 \pi y\left|\xi^{(1)}\right|\right) K_{s_{2}-\frac{1}{2}}\left(2 \pi y\left|\xi^{(2)}\right|\right) y^{s_{0}-1} \frac{\mathrm{~d} y}{y}
\end{aligned}
$$

## More Precisely II

## Proposition

## where

$$
\begin{aligned}
& V_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)=\zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{2}-1\right)\left(t_{1} t_{2}\right)^{-\frac{3 \nu_{1}-3 \nu_{2}+3}{2}} \\
& E^{\star}\left(\frac{t_{1} \alpha_{1}^{(1)}-\mathrm{i} t_{2} \alpha_{1}^{(2)}}{t_{1} \alpha_{2}^{(1)}-\mathrm{i} t_{2} \alpha_{2}^{(2)}}, \frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right)
\end{aligned}
$$

and $\left(\alpha_{1}, \alpha_{2}\right)$ is a $\mathbb{Z}$-basis of $\mathfrak{a}$.

## What about $\Omega_{5}$ and $\Omega_{6}$ ?

## It can be shown that



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$G_{\nu_{1}, \nu_{2}}^{\Omega_{6}}\left(t_{1}, t_{2}\right)=\left(t_{1} t_{2}\right)^{-3 \nu_{2}-\frac{3 \nu_{1}}{2}} \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right)$


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$$

$$
\times E^{\star}\left(\frac{t_{1} \alpha_{1}^{(1)}+\mathrm{i} t_{2} \alpha_{1}^{(2)}}{t_{1} \alpha_{2}^{(1)}+\mathrm{i} t_{2} \alpha_{2}^{(2)}}, \frac{3 \nu_{2}}{2}\right)
$$

$G_{\nu_{1}, \nu_{2}}^{\Omega_{6}}\left(t_{1}, t_{2}\right)=\left(t_{1} t_{2}\right)^{-3 \nu_{2}-\frac{3 \nu_{1}}{2}} \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right)$

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$$
\begin{aligned}
& G_{\nu_{1}, \nu_{2}}^{\Omega_{5}}\left(t_{1}, t_{2}\right)=\left(t_{1} t_{2}\right)^{3 \nu_{1}+\frac{3 \nu_{2}}{2}} \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}\right.\left.+3 \nu_{2}-1\right) \\
& \times E^{\star}\left(\frac{t_{1} \alpha_{1}^{(1)}+\mathrm{i} t_{2} \alpha_{1}^{(2)}}{t_{1} \alpha_{2}^{(1)}+\mathrm{i} t_{2} \alpha_{2}^{(2)}}, \frac{3 \nu_{2}}{2}\right) \\
&\left.G_{\nu_{1}, \nu_{2}}^{\Omega_{6}}\left(t_{1}, t_{2}\right)=\left(t_{1} t_{2}\right)^{-3 \nu_{2}-\frac{3 \nu_{1}}{2} \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}\right.}+3 \nu_{2}-1\right) \\
& \times E^{\star}\left(\frac{t_{1} \alpha_{1}^{(1)}+\mathrm{i} t_{2} \alpha_{1}^{(2)}}{t_{1} \alpha_{2}^{(1)}+\mathrm{i} t_{2} \alpha_{2}^{(2)}}, \frac{3 \nu_{1}}{2}\right)
\end{aligned}
$$

## In Summary

## Letting


we obtain

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$$
f_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)=
$$

$$
\zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{2}-1\right)\left(t_{1} t_{2}\right)^{-\frac{3 \nu_{1}-3 \nu_{2}+3}{2}} E^{\star}\left(\frac{t_{1} \alpha_{1}^{(1)}-\mathrm{i} t_{2} \alpha_{1}^{(2)}}{t_{1} \alpha_{2}^{(1)}-\mathrm{i} t_{2} \alpha_{2}^{(2)}}, \frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right)
$$



## we obtain

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$$
\begin{aligned}
& f_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)= \\
& \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{2}-1\right)\left(t_{1} t_{2}\right)^{-\frac{3 \nu_{1}-3 \nu_{2}+3}{2}} E^{\star}\left(\frac{t_{1} \alpha_{1}^{(1)}-\mathrm{i} t_{2} \alpha_{1}^{(2)}}{t_{1} \alpha_{2}^{(1)}-\mathrm{i}_{2} \alpha_{2}^{(2)}}, \frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \\
& \quad+\left(t_{1} t_{2}\right)^{3 \nu_{1}+\frac{3 \nu_{2}}{2} \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{1}-1\right) E^{\star}\left(\frac{t_{1} \alpha_{1}^{(1)}+\mathrm{i} t_{2} \alpha_{1}^{(2)}}{t_{1} \alpha_{2}^{(1)}+\mathrm{i} t_{2} \alpha_{2}^{(2)}}, \frac{3 \nu_{2}}{2}\right)} \\
& \quad+\left(t_{1} t_{2}\right)^{-3 \nu_{2}-\frac{3 u_{1}}{2} \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) E^{\star}\left(\frac{t_{1} \alpha_{1}^{(1)}+\mathrm{i} t_{2} \alpha_{1}^{(2)}}{t_{1} \alpha_{2}^{(1)}+\mathrm{i} t_{2} \alpha_{2}^{(2)}}, \frac{3 \nu_{1}}{2}\right)}
\end{aligned}
$$

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\begin{aligned}
& f_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)= \\
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& \quad+\left(t_{1} t_{2}\right)^{3 \nu_{1}+\frac{3 \nu_{2}}{2}} \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{1}-1\right) E^{\star}\left(\frac{t_{1} \alpha_{1}^{(1)}+\mathrm{i} t_{2} \alpha_{1}^{(2)}}{t_{1} \alpha_{2}^{(1)}+\mathrm{i} t_{2} \alpha_{2}^{(2)}}, \frac{3 \nu_{2}}{2}\right) \\
& \quad+\left(t_{1} t_{2}\right)^{-3 \nu_{2}-\frac{3 \nu_{1}}{2} \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) E^{\star}\left(\frac{t_{1} \alpha_{1}^{(1)}+\mathrm{i} t_{2} \alpha_{1}^{(2)}}{t_{1} \alpha_{2}^{(1)}+\mathrm{i} t_{2} \alpha_{2}^{(2)}}, \frac{3 \nu_{1}}{2}\right)}
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& \quad+\left(t_{1} t_{2}\right)^{3 \nu_{1}+\frac{3 \nu_{2}}{2}} \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{1}-1\right) E^{\star}\left(\frac{t_{1} \alpha_{1}^{(1)}+\mathrm{i} t_{2} \alpha_{1}^{(2)}}{t_{1} \alpha_{2}^{(1)}+\mathrm{i} t_{2} \alpha_{2}^{(2)}}, \frac{3 \nu_{2}}{2}\right) \\
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\end{aligned}
$$

we obtain

## In Summary II

## Theorem

Same hypotheses as before. We have

$$
R_{A}\left(s_{0}, s_{1}, s_{2}\right)=6 \iint_{\left(\mathbb{R}_{+}^{*}\right)^{2} / \mathfrak{o}^{\times}}\left(G_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)-f_{\nu_{1}, \nu_{2}}\left(t_{1}, t_{2}\right)\right)\left(t_{1} t_{2}\right)^{s} \frac{d t_{1} d t_{2}}{t_{1} t_{2}}
$$

## What have we got so far?

- So, we have an identity linking our renormalized integral $R_{A}$, $G_{\nu_{1}, \nu_{2}}$ and $f_{\nu_{1}, \nu_{2}}$.
- Although $w$ is a functional equation for $G_{\nu_{1}, \nu_{2}}$, it does not leave $f_{\nu_{1}, \nu_{2}}$ invariant.
- But, as it turns out, the integration over $\left(\mathbb{R}_{+}^{*}\right)^{2} / 0^{x}$ has the effect of creating this extra functional equation.
- This requires some work, though.


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Introduction and First Notations
Polar divisor considerations
Proof of the Main Theorem Acknowledgements

## Hmm... Kinda nasty...



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$$
\frac{R_{A}\left(s_{0}, s_{1}, s_{2}\right)}{3}=\int_{0}^{+\infty}\left(\mathcal{G}_{\nu_{1}, \nu_{2}}(v)\right.
$$

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$$
\begin{aligned}
\frac{R_{A}\left(s_{0}, s_{1}, s_{2}\right)}{3}= & \int_{0}^{+\infty}\left(\mathcal{G}_{\nu_{1}, \nu_{2}}(v)\right. \\
& -v^{-\frac{3 \nu_{1}-3 \nu_{2}+3}{2}} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{2}-1\right) \\
& -v^{3 \nu_{1}+\frac{3 \nu_{2}}{2} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right)} \\
& -v^{\left.-3 \nu_{2}-\frac{3 \nu_{1}}{2} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right)\right) v^{s} \frac{d v}{v}}
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& \quad-v^{3 \nu_{1}+\frac{3 \nu_{2}}{2}} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right)
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& \left.-v^{-3 \nu_{2}-\frac{3 \nu_{1}}{2}} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right)\right) v^{s} \frac{\mathrm{~d} v}{v}
\end{aligned}
$$

## Final Modifications

- Replace $v$ by $\frac{1}{v}$ in the first one.
- Regroup into one integral from 1 to
- Let $x=\sqrt{v}$, in order to get rid of those fractional powers.
- Shake.


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## Nastier

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\begin{aligned}
\frac{R_{A}\left(s_{0}, s_{1}, s_{2}\right)}{3}= & \int_{1}^{+\infty}\left(\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{2}\right) x^{2 s}+\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{-2}\right) x^{-2 s}\right. \\
& -x^{-\left(3 \nu_{1}+3 \nu_{2}+3-2 s\right)} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{2}-1\right) \\
& -x^{6 \nu_{1}+3 \nu_{2}+2 s} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
& -x^{-\left(3 \nu+6 \nu_{2}-2 s\right)} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
& -x^{3 \nu_{1}-3 \nu_{2}+3-2 s} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}-2\right) \\
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& \left.-x^{3 \nu_{1}+6 \nu_{2}-2 s} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right)\right) \frac{\mathrm{d} x}{x}
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& -\frac{1}{3 \nu_{1}-3 \nu_{2}+3-2 s} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{2}-1\right) \\
& -\frac{1}{3 \nu_{1}+6 \nu_{2}-2 s} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
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& \left.-x^{3 \nu_{1}+6 \nu_{2}-2 s} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right)\right) \frac{\mathrm{d} x}{x} \\
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& -\frac{1}{6 \nu_{1}+3 \nu_{2}+2 s} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right)
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& -x^{-\left(3 \nu_{1}+6 \nu_{2}-6-2 s\right)} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}-1\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \\
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\end{aligned}
$$

## Even More Nastier II

$$
\begin{aligned}
& -\frac{1}{3 \nu_{1}-3 \nu_{2}+3-2 s} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{2}-1\right) \\
& -\frac{1}{3 \nu_{1}+6 \nu_{2}-2 s} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
& -\frac{1}{6 \nu_{1}+3 \nu_{2}+2 s} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
& +\frac{1}{3 \nu_{1}+6 \nu_{2}-6-2 s} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}-1\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \\
& +\frac{1}{3 \nu_{1}-3 \nu_{2}-3-2 s} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}-1\right) \zeta^{\star}\left(3 \nu_{2}\right) \\
& +\frac{1}{6 \nu_{1}+3 \nu_{2}+2 s-6} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \zeta^{\star}\left(3 \nu_{1}-1\right) .
\end{aligned}
$$

## Even More Nastier

$$
\begin{aligned}
\frac{R_{A}\left(s_{0}, s_{1}, s_{2}\right)}{3}= & \int_{1}^{+\infty}\left(\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{2}\right) x^{2 s}+\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{-2}\right) x^{-2 s}\right. \\
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\quad & \left.-x^{-\left(6 \nu_{1}+3 \nu_{2}-6+2 s\right)} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \zeta^{\star}\left(3 \nu_{1}-1\right)\right) \frac{\mathrm{d} x}{x}
\end{aligned}
$$

## Even More Nastier

$$
\begin{aligned}
\frac{R_{A}\left(s_{0}, s_{1}, s_{2}\right)}{3}= & \int_{1}^{+\infty}\left(\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{2}\right) x^{2 s}+\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{-2}\right) x^{-2 s}\right. \\
\square & -x^{6 \nu_{1}+3 \nu_{2}+2 s} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
\square & -x^{3 \nu_{1}-3 \nu_{2}+3-2 s} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}-2\right) \\
\quad & -x^{3 \nu_{1}+6 \nu_{2}-2 s} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
\square & -x^{-\left(3 \nu_{1}+6 \nu_{2}-6-2 s\right)} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}-1\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \\
& -x^{-\left(3 \nu_{1}-3 \nu_{2}-3-2 s\right)} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}-1\right) \zeta^{\star}\left(3 \nu_{2}\right) \\
\quad & \left.-x^{-\left(6 \nu_{1}+3 \nu_{2}-6+2 s\right)} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \zeta^{\star}\left(3 \nu_{1}-1\right)\right) \frac{\mathrm{d} x}{x}
\end{aligned}
$$

## Even More Nastier

$$
\begin{aligned}
\frac{R_{A}\left(s_{0}, s_{1}, s_{2}\right)}{3}= & \int_{1}^{+\infty}\left(\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{2}\right) x^{2 s}+\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{-2}\right) x^{-2 s}\right. \\
\square & -x^{6 \nu_{1}+3 \nu_{2}+2 s} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
\square & -x^{3 \nu_{1}-3 \nu_{2}+3-2 s} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}-2\right) \\
\square & -x^{3 \nu_{1}+6 \nu_{2}-2 s} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
\square & -x^{-\left(3 \nu_{1}+6 \nu_{2}-6-2 s\right)} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}-1\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \\
\square & -x^{-\left(3 \nu_{1}-3 \nu_{2}-3-2 s\right)} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}-1\right) \zeta^{\star}\left(3 \nu_{2}\right) \\
\square & \left.-x^{-\left(6 \nu_{1}+3 \nu_{2}-6+2 s\right)} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \zeta^{\star}\left(3 \nu_{1}-1\right)\right) \frac{\mathrm{d} x}{x}
\end{aligned}
$$

## Even More Nastier

$$
\begin{aligned}
\frac{R_{A}\left(s_{0}, s_{1}, s_{2}\right)}{3}= & \int_{1}^{+\infty}\left(\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{2}\right) x^{2 s}+\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{-2}\right) x^{-2 s}\right. \\
\square & -x^{6 \nu_{1}+3 \nu_{2}+2 s} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
\square & -x^{3 \nu_{1}-3 \nu_{2}+3-2 s} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}-2\right) \\
\square & -x^{3 \nu_{1}+6 \nu_{2}-2 s} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
\square & -x^{-\left(3 \nu_{1}+6 \nu_{2}-6-2 s\right)} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}-1\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \\
\square & -x^{-\left(3 \nu_{1}-3 \nu_{2}-3-2 s\right)} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}-1\right) \zeta^{\star}\left(3 \nu_{2}\right) \longrightarrow \\
\longrightarrow & \left.-x^{-\left(6 \nu_{1}+3 \nu_{2}-6+2 s\right)} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \zeta^{\star}\left(3 \nu_{1}-1\right)\right) \frac{\mathrm{d} x}{x}
\end{aligned}
$$

## Even More Nastier

$$
\begin{aligned}
\frac{R_{A}\left(s_{0}, s_{1}, s_{2}\right)}{3}= & \int_{1}^{+\infty}\left(\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{2}\right) x^{2 s}+\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{-2}\right) x^{-2 s}\right. \\
\square & -x^{6 \nu_{1}+3 \nu_{2}+2 s} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
\square & -x^{3 \nu_{1}-3 \nu_{2}+3-2 s} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}-2\right) \\
\quad & -x^{3 \nu_{1}+6 \nu_{2}-2 s} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
& -x^{-\left(3 \nu_{1}+6 \nu_{2}-6-2 s\right)} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}-1\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \\
\square & -x^{-\left(3 \nu_{1}-3 \nu_{2}-3-2 s\right)} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}-1\right) \zeta^{\star}\left(3 \nu_{2}\right) \longrightarrow \\
\longrightarrow & \left.x^{-\left(6 \nu_{1}+3 \nu_{2}-6+2 s\right)} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \zeta^{\star}\left(3 \nu_{1}-1\right)\right) \frac{\mathrm{d} x}{x}
\end{aligned}
$$

## Even More Nastier

$$
\begin{aligned}
& \frac{R_{A}\left(s_{0}, s_{1}, s_{2}\right)}{3}=\int_{1}^{+\infty}\left(\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{2}\right) x^{2 s}+\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{-2}\right) x^{-2 s}\right. \\
& \square-x^{6 \nu_{1}+3 \nu_{2}+2 s} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
& \text { - }-x^{3 \nu_{1}-3 \nu_{2}+3-2 s} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}-2\right) \\
& \longrightarrow-x^{3 \nu_{1}+6 \nu_{2}-2 s} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
& \longrightarrow-x^{-\left(3 \nu_{1}+6 \nu_{2}-6-2 s\right)} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}-1\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \\
& -x^{-\left(3 \nu_{1}-3 \nu_{2}-3-2 s\right)} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}-1\right) \zeta^{\star}\left(3 \nu_{2}\right) \\
& \left.\longrightarrow-x^{-\left(6 \nu_{1}+3 \nu_{2}-6+2 s\right)} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \zeta^{\star}\left(3 \nu_{1}-1\right)\right) \frac{\mathrm{d} x}{x}
\end{aligned}
$$

## Even More Nastier

$$
\begin{aligned}
& \frac{R_{A}\left(s_{0}, s_{1}, s_{2}\right)}{3}=\int_{1}^{+\infty}\left(\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{2}\right) x^{2 s}+\mathcal{G}_{\nu_{1}, \nu_{2}}\left(x^{-2}\right) x^{-2 s}\right. \\
& \text { - }-x^{6 \nu_{1}+3 \nu_{2}+2 s} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
& \text { - }-x^{3 \nu_{1}-3 \nu_{2}+3-2 s} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}\right) \zeta^{\star}\left(3 \nu_{1}-2\right) \\
& \longrightarrow-x^{3 \nu_{1}+6 \nu_{2}-2 s} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-1\right) \\
& { }^{-\left(3 \nu_{1}+6 \nu_{2}-6-2 s\right)} L\left(\frac{3 \nu_{1}}{2}\right) \zeta^{\star}\left(3 \nu_{2}-1\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \\
& -x^{-\left(3 \nu_{1}-3 \nu_{2}-3-2 s\right)} L\left(\frac{3 \nu_{1}+3 \nu_{2}-1}{2}\right) \zeta^{\star}\left(3 \nu_{1}-1\right) \zeta^{\star}\left(3 \nu_{2}\right) \\
& \left.\longrightarrow-x^{-\left(6 \nu_{1}+3 \nu_{2}-6+2 s\right)} L\left(\frac{3 \nu_{2}}{2}\right) \zeta^{\star}\left(3 \nu_{1}+3 \nu_{2}-2\right) \zeta^{\star}\left(3 \nu_{1}-1\right)\right) \frac{\mathrm{d} x}{x}
\end{aligned}
$$

## Even More Nastier II



## Acknowledgements

## Thanks, yall!

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